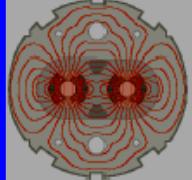
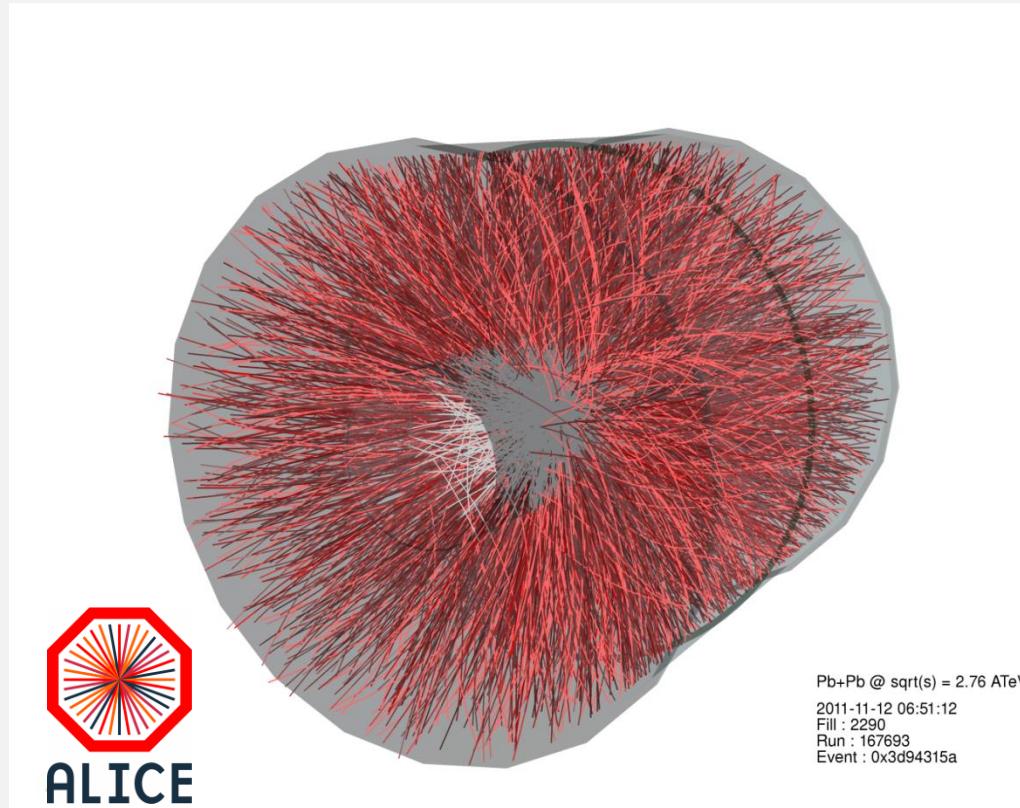


Fyzika s t'ažkými iónmi

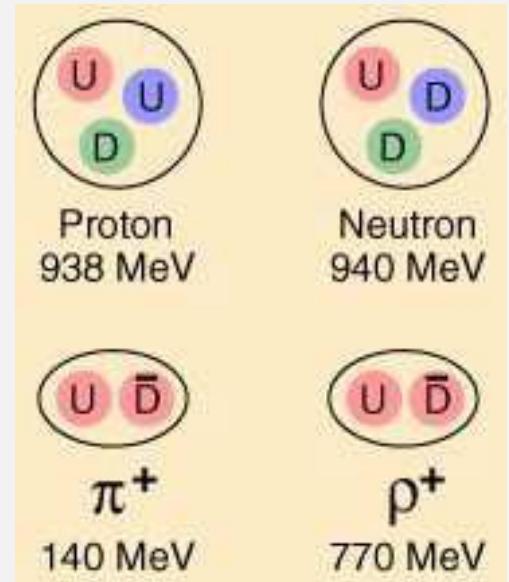
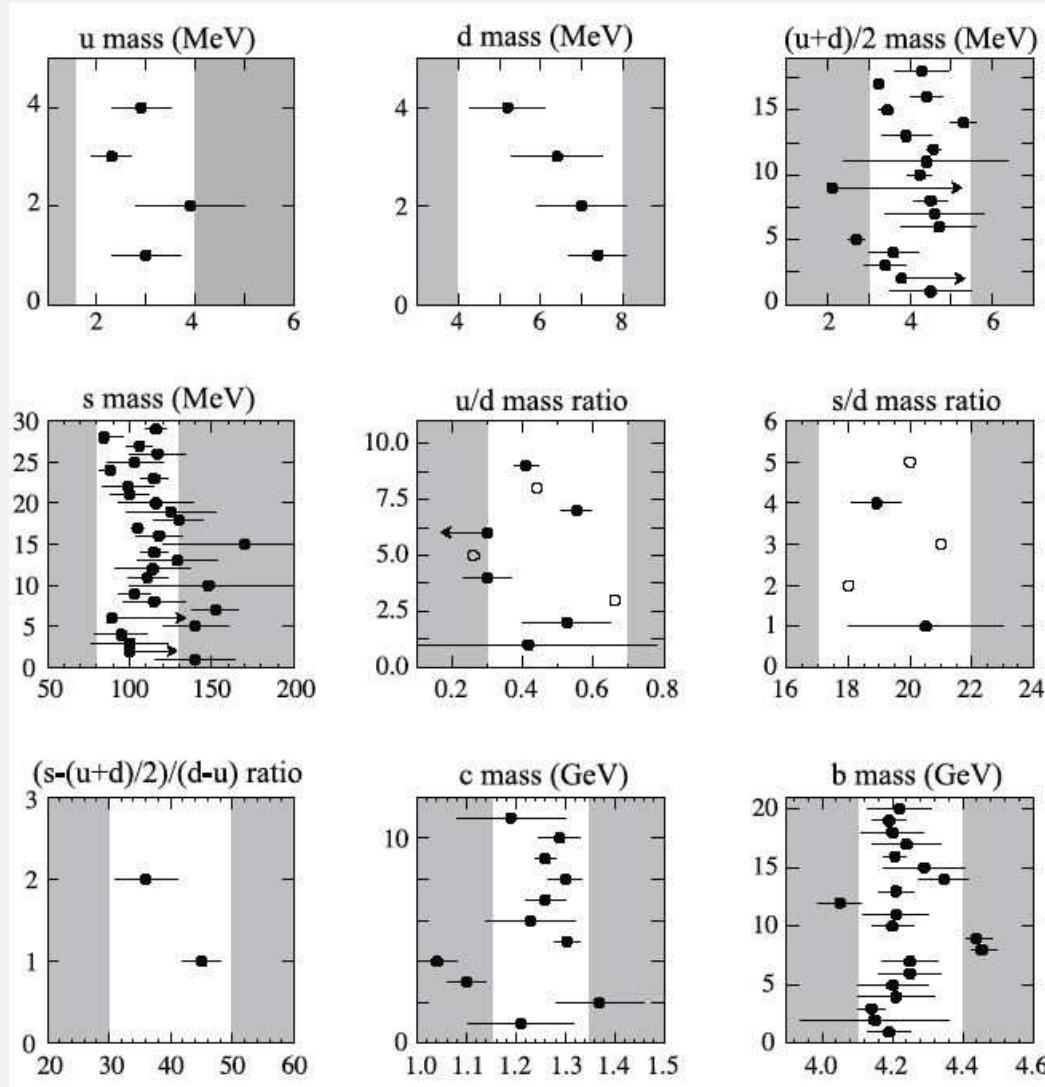
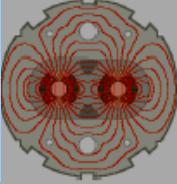


alebo malý tresk



Karel Šafařík, CERN

Hmotnosť nukleónov

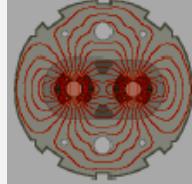


Σ hmotnosti kvarkov ~
~ 1 % hmotnosti p/n !

väzbová energia

Vákum v QED a QCD

kvantová elektrodynamika: QED
kvantová chromodynamika: QCD



- ◆ **Energia páru nábojov spontánne narodených vo vákuu – kvantová fluktuácia ($\hbar = 1$, $c = 1$):**

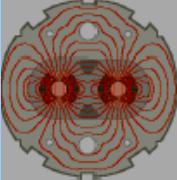
$$E_{\text{kin}} = p \sim 1/r \quad (p \times r \geq 1)$$

$$E_{\text{pot}} = -q^2/(4\pi r) \quad (q = e \text{ or } q = g_s)$$

$$E = E_{\text{kin}} + E_{\text{pot}} = (1/r) \times (1 - q^2/4\pi)$$

- ◆ **v QED toto je pravda pre lubovlonú “škálu” (už po Planckovu “škálu” $\sim 10^{-20}$ fm)**
- ◆ **v QCD to je však správne len pre velmi malé vzdialenosť, niekol'ko fm (10^{-13} cm)**

Prípad QED



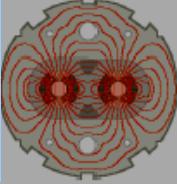
◆ v QED

$$q^2 = e^2 = 4\pi\alpha_{em}$$

- α_{em} sa mení zo vzdialenosťou (polarizácia vákua)
- kde pre veľké vzdialosti $\alpha_{em} = 1/137$
- pri EW (elektro-slabej) škále ($r = 2 \times 10^{-3}$ fm) $\alpha_{em} = 1/128$
- pri Planckovej “škále” ($r = 10^{-20}$ fm) $\alpha_{em} = 1/76$

- ## ◆ To znamená, že číselný faktor pred $1/r$: $(1 - q^2/4\pi)$ sa mení zo vzdialosťou, ale
- len málo, medzi 0.987 – 0.993 (i.e. 0.6%) ak meníme vzdialenosť od Planckovej “škály” až po nekonečno...

Prípad QCD



◆ v QCD

$$q^2 = g_s^2 = 4\pi\alpha_s$$

- kde α_s sa zmenšuje veľmi rýchlo zo vzdialostou (asymptotická sloboda)

- pri Planckovej škále $\alpha_s = 0.04$
- pri elektro-slabej skále $\alpha_s = 0.118$
- pri $\Lambda_{QCD} \approx 0.2 \text{ GeV}$ ($r \approx 1 \text{ fm}$) $\alpha_s \approx 1$

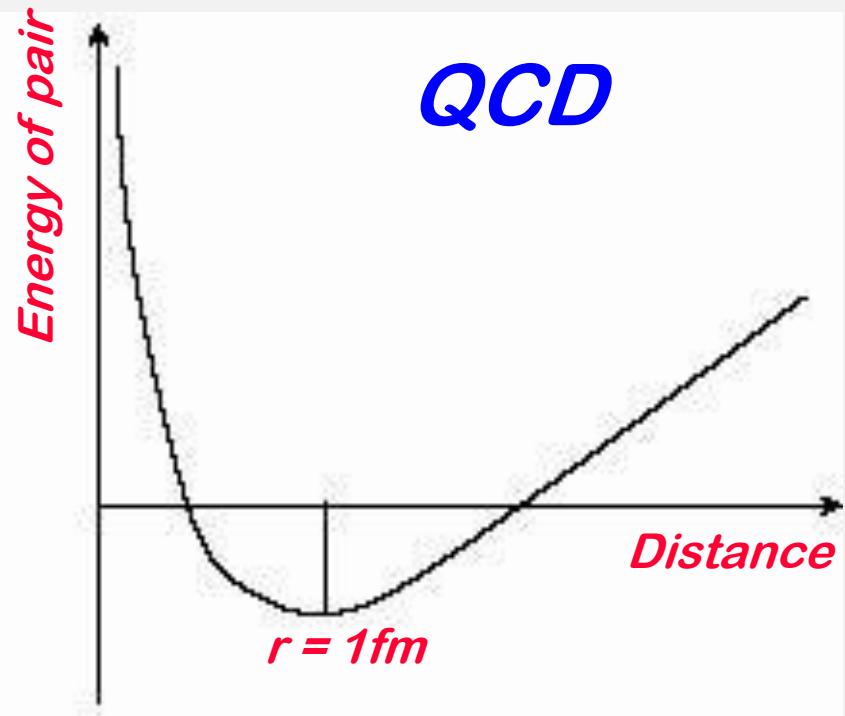
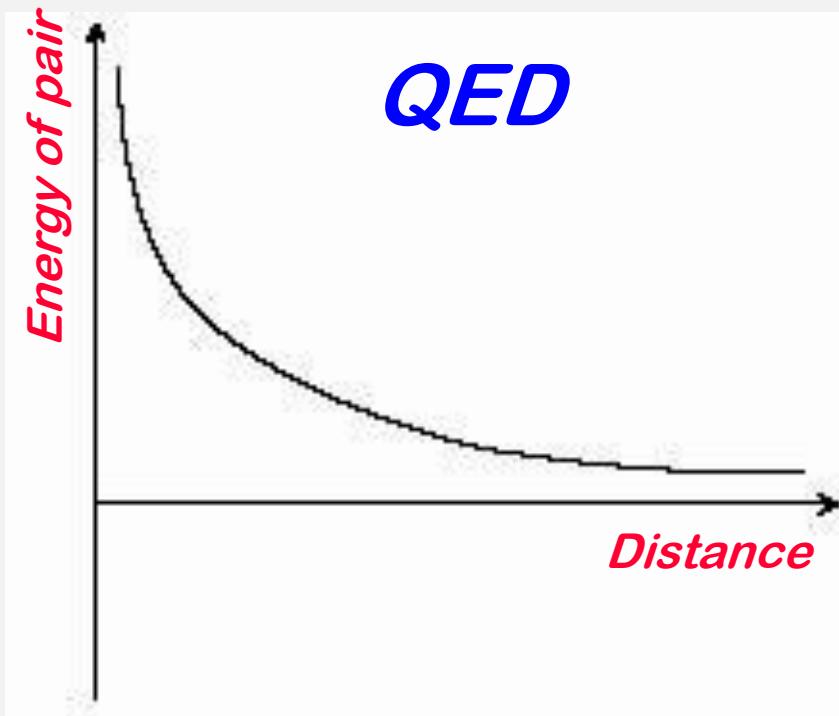
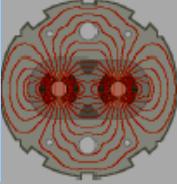
◆ numerický faktor $(1 - q^2/4\pi) = 1 - \alpha_s$

- sa znižuje so vzdialenosťou, pri Planckovej škále je 0.96
- ale pozor, pre $r \approx 1 \text{ fm}$ uz je záporný !

◆ pri väčších vzdialenosťach je $E = \sigma \times r$ ($\sigma \approx 1 \text{ GeV/fm}$)

- a tento faktor je opäť kladný

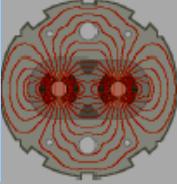
QED versus QCD



Kinetická energia stále dominuje nad potenciálnou (pole je slabé) virtual páry

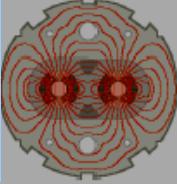
Energia skrytá v poli prevázi pri nejakej vzdialosti kinetickú reálne páry – vakuový kondensát

Symetrie QCD



- ◆ **QCD má dve približné symetrie:**
 - Z_3 –(centre) symetriu (pre čisto kalibračnú toer, v limite $m_q \rightarrow \infty$)
 - chirálnu symetriu (obnovenú pre nulové hmotnosti, t.j. $m_q \rightarrow 0$)
- ◆ **Pri vel'kých hustotách a teplotách sa nakoniec**
 - Z_3 –symetria naruší (prechod od confinementu k deconfinementu)
 - chirálna symetria obnoví (chirálny fázový prechod)
- ◆ **Otázky:**
 - existuje jeden spoločný fázový prechod alebo dva nezávislé?
 - akého druhu je tento (tieto) fázový(é) prechod(y)
 - prvého druhu (má latentné teplo) ?
 - druhého druhu (len zlom) ?
 - alebo je to len cross-over prechod ?

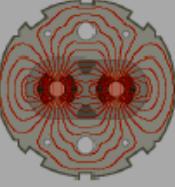
Chirálna symetria



◆ Pre $m_q \rightarrow 0$ helicita kavarkov sa zachováva

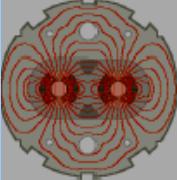
- pretože gluóny majú helicity ± 1 QCD teória v tejto limite má $SU(3)_L \times SU(3)_R$ symetriu
 - QCD svet sa rozpadol na dva svety ktoré navzájom nekomunikujú – lavácky svet a pravácky
- ak dáme do QCD vákuu nehmotný lavotočivý kvark, on môže anihilovať s lavotočivým anti-kvarkom z vákuového kondenzátu – tým sa ale oslobodí pravotočivý kvark
 - pre vzdialeného pozorovatela naš testový kvark spontánne zmenil helicitu a preto musel nejako získať dynamickú hmotnosť !
 - QCD kvark—anti-kvarkový kondenzát generuje dynamické kvarkové hmotnosť a narušuje chirálnu symetriu
- ak zvýšime teplotu kinetická energia nabitého páru (nad nejakou hodnotou) prevýši potenciálnu energiu
 - kvark—anti-kvark kondenzát zmizne z vákuu
 - chirálna symetria sa obnoví nad nejakou kritickou teplotou
 - hodnota $\langle 0 | \bar{q}q | 0 \rangle$ je “order parameter” fázového prechodu

Confinement (uveznenie)



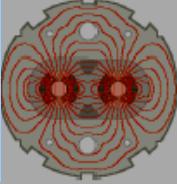
- ◆ **hmotné kvarky v čisto gluónovom vákuume pri nulovej teplote**
 - nie sú viditeľné detektorom kvôli deštruktívnej interferencii
 - expectation hodnota pre stopu kvarkového propagátoru – 3– hodnotový path integrál s rôznymi fázami
$$\exp(i \times 2\pi j/3), \quad j=1,2,3 \quad (\text{generátory } Z_3)$$
 - zvyšujúc teploty T až po nejakú hodnotu toto zostane tak
 - až pokial gluónové pole bude mať dostatok času sledovať' (koherentný rearangement) náš testový farebný náboj
 - Dalšie zvýšenie teploty (nad nejakú kritickú hodnotu) gluónové pole nebude mať dostatok času
 - Interferencia troch ciest sa naruší
 - test farebný náboj sa stane detektovateľný, bude deconfinovaný
 - Toto sa dá spočítať analytickým predĺžením kvarkového propagátoru v komplexnom case ($t = +i/T$) – Polyakov loop – ktorý sa stane nenulovým pre $T > T_c$
 - Polyakov loop je “order parameter” fázového prechodu

Symetrie QCD



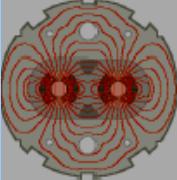
- ◆ **Obidve symetrie sú narušené dynamicky**
 - Z_3 symmetria je narušená kinetickou energiou (pri vysokej T)
 - order parameter (Polyakov loop) je nulový pod T_c a nenulový nad
 - je to “order – disorder” fázový prechod, Z_3 je narušená nad T_c
 - chirálna symetria je narušená potenciálnou energiou (pri nízkej T)
 - order parameter (kvark—anti-kvarkový kondenzát) je nenulový pod T_c a nulový nad
 - je to “disorder – order” fázový prechod, chirálna symetria je obnovená nad T_c
- ◆ **Obidve sú však narušené aj explicitne – hmotnosťou**
 - pre malosť m_q je reálne že scénar ohľadom chirálnej symetrie zostane dobrým priblížením
 - ale čo so Z_3 symetriou, prečo nie je úplne zničená malosťou m_q ?

Obnovenie konfinementu



- ◆ **Ked sa snažíme znižiť m_q z nekonečna na ich vlastnú (malú) hodnotu to co sa stane závisí od teploty:**
 - pri nízkych teplotách m_q sa efektívne prestane znižovať ked' prídeme pod dynamickú hmotnosť kvarku $M_q \approx 350 \text{ MeV}$ pretože chirálna symetria je narušená
 - Z_3 symetria zostane približnou symetriou pri nízkych teplotách aj po takomto tvrdom pokuse o explicitné narušenie
 - narušenie chirálnej symetrie efektívne zvyšuje hmotnosti kvarkov a preto riadi obnovenie Z_3 symetrie
 - toto je argument preto, aby obidva fázové prechody nastali v tom istom bode

Fázy QCD – hračkársky model



◆ uvážme

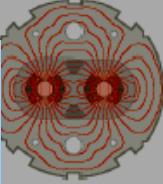
- fázu v confinemente (hadrónový plyn, HG) z piónov
- deconfinovanú fazu (kvark—gluónovú plazmu, QGP) z gluónov a dvoch typov (vôní) kvarkov
- stavové rovnica pre ideálny plyn

$$\varepsilon = (g/30) \pi^2 T^4, \quad p = \varepsilon/3 = (g/90) \pi^2 T^4$$

$$\text{kde } g = n_b + (7/8) n_f$$

- pre HG $n_b = 3, n_f = 0$
 $p_{HG} = (1/30) \pi^2 T^4$
- pre QGP: $n_b = 16, n_f = 24$ ale teraz máme aj vonkajší tlak od QCD vákua B
 $p_{QGP} = (37/90) \pi^2 T^4 - B$
- na hranici dvoch fáz – tlak musí byť rovnaký
 $T_c = (90B/34\pi^2)^{1/4} = 144 \text{ MeV}$ pre $B^{1/4} = 200 \text{ MeV}$ (**MIT bag model**)

Fázy QCD – poruchová teória



- ◆ **pri nenulovej baryónovej hustote – prvý rád p-QCD**

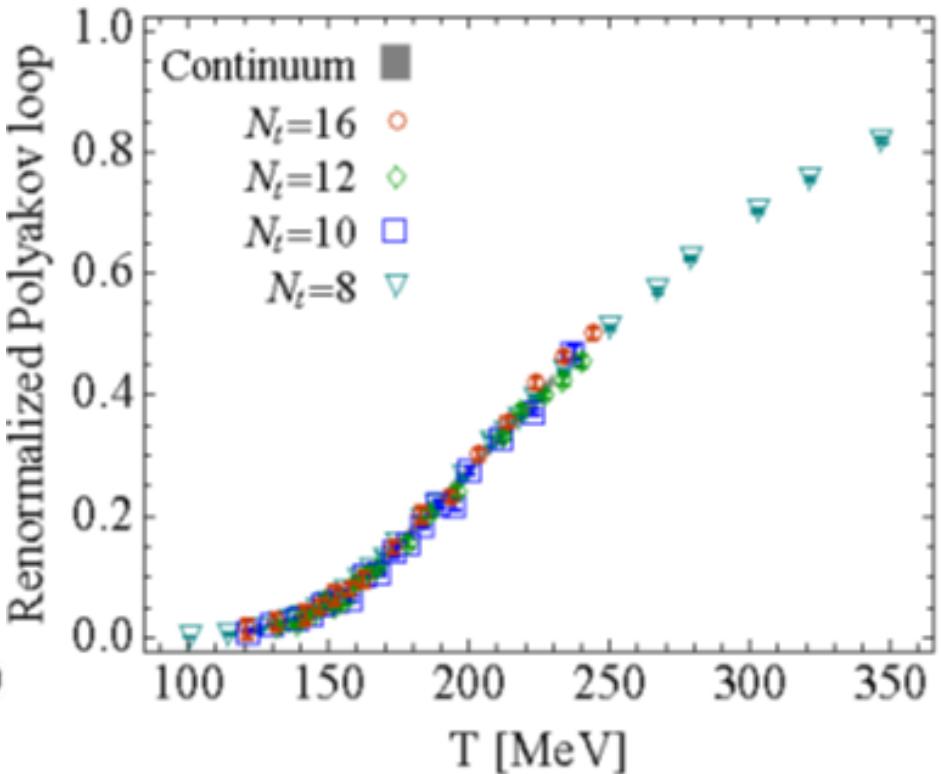
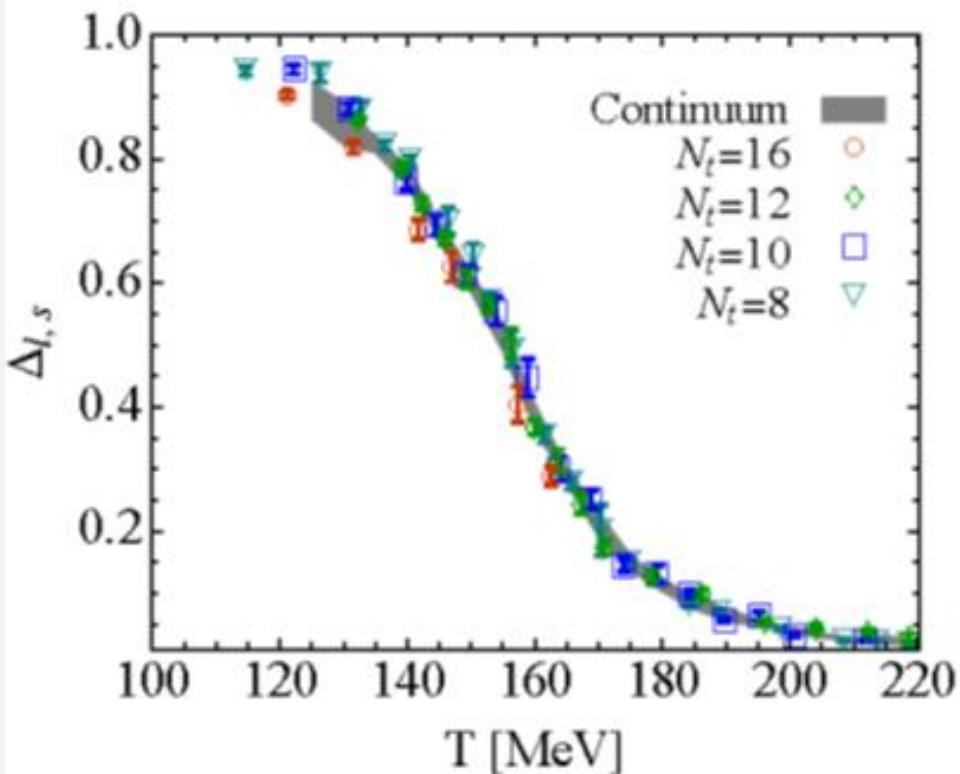
$$\begin{aligned}\varepsilon = & [16(1 - 15\alpha_s/4\pi) + (7/8)12n_q(1 - 50\alpha_s/21\pi)] (1/30) \pi^2 T^4 + \\ & + \sum_q 16(1 - 15\alpha_s/2\pi) (3/\pi^2) \mu_q^{-2} (\pi^2 T^4 + \mu_q^{-2}) / 2\end{aligned}$$

(pre $\mu_q = 0$, $\alpha_s = 0$, a $n_q = 2$ dostaneme náš hračkársky model)

použijúc $\alpha_s = 0.4$ tou istou cestou dostaneme $T_c = 164$ MeV

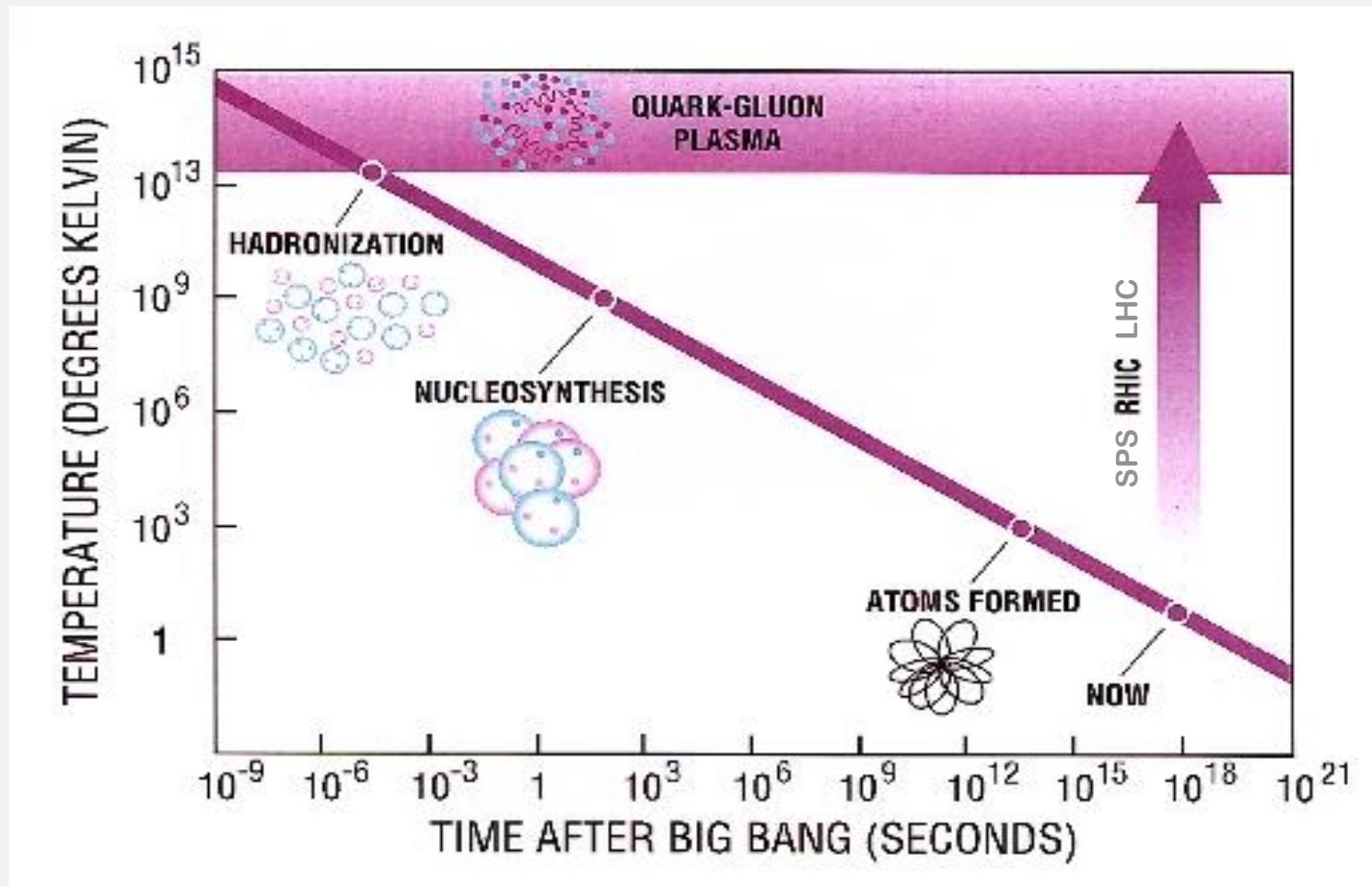
- ◆ Dnes analytické výpočty existujú aj pre vyššie rády

Lattice QCD at high temperature

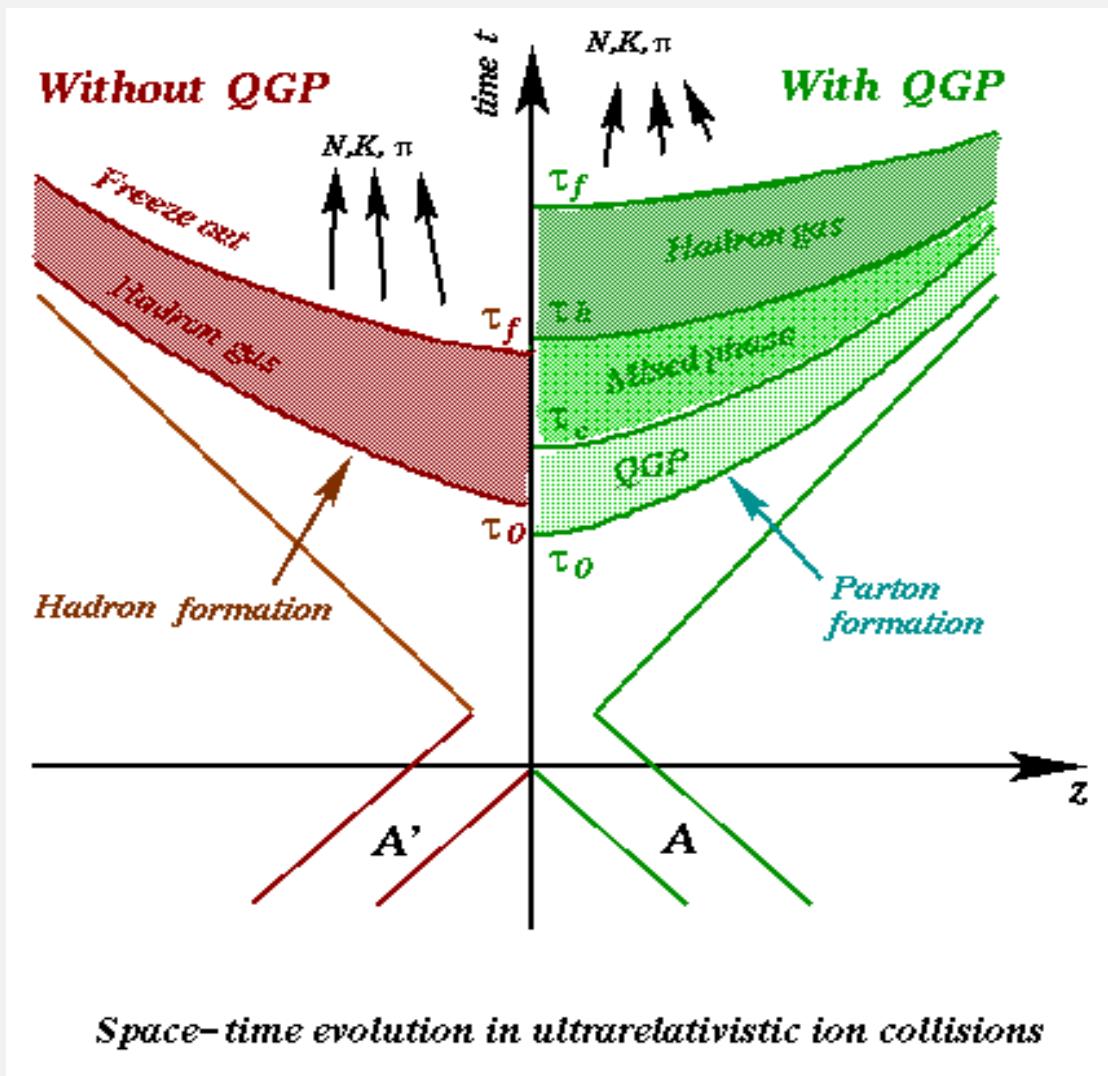


- Order parameters of QCD phase transitions:
 - quark – antiquark vacuum condensate
 - chiral phase transition in limit $m_q \rightarrow 0$
 - expectation value of Polyakov loop
 - deconfinement phase transition in limit $m_q \rightarrow \infty$

Big Bang



Priestorovo-časový vývoj



R_{AA} and v_n – definitions

- R_{AA} – ratio of p_T spectrum in AA collisions to that in pp
– properly normalized by number of binary collisions

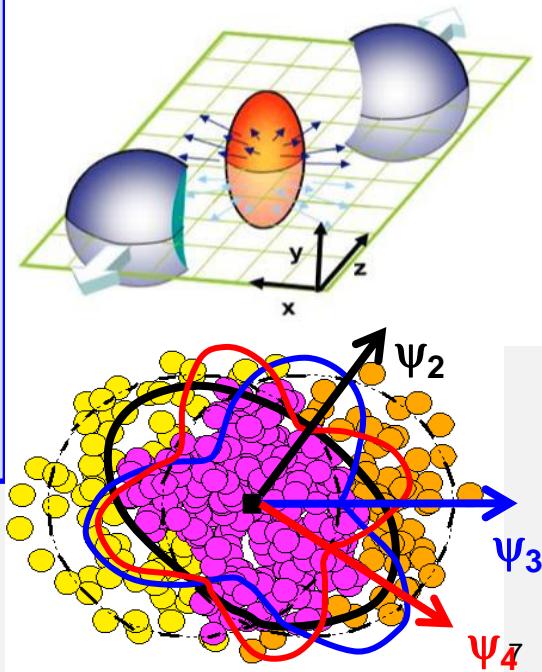
$$R_{AA} = \frac{(d\sigma/dp_T)_{AA}}{N_{\text{bin}}(d\sigma/dp_T)_{pp}} = \dots$$

if AA would be just a superposition of pp collisions $R_{AA} = 1$

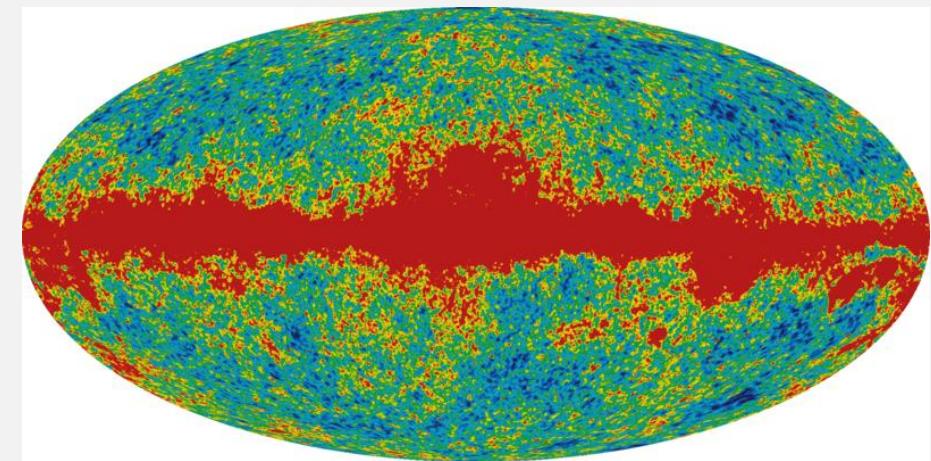
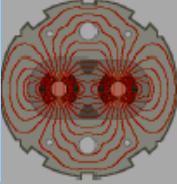
- v_n – Fourier coefficients of particle distributions in azimuthal angle φ with respect to n -th reaction plane

$$\frac{dN}{d\varphi}(\dots) \propto 1 + 2 \sum_{n=1}^{\infty} v_n \cos n(\varphi - \psi_n)$$

$v_n = 0$ would mean azimuthally symmetric distribution

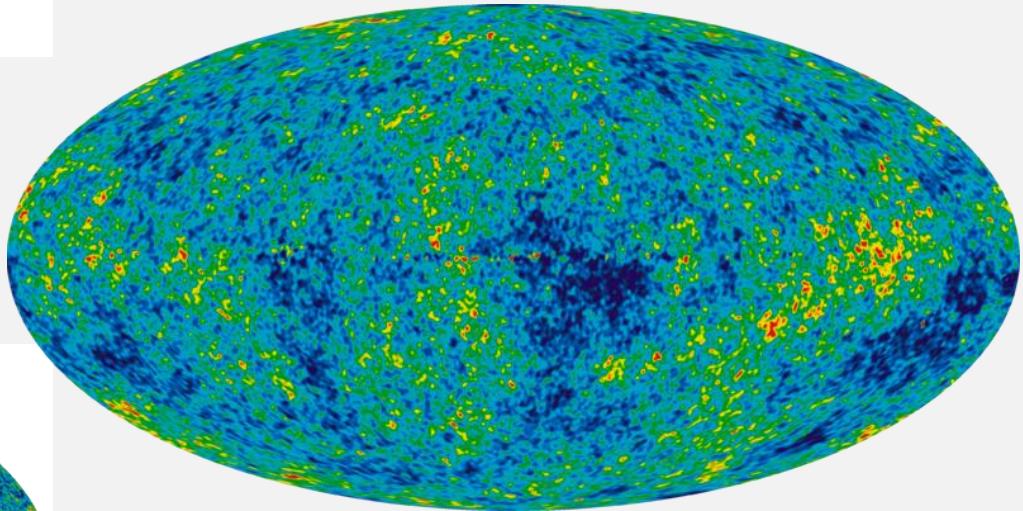


CMB maps

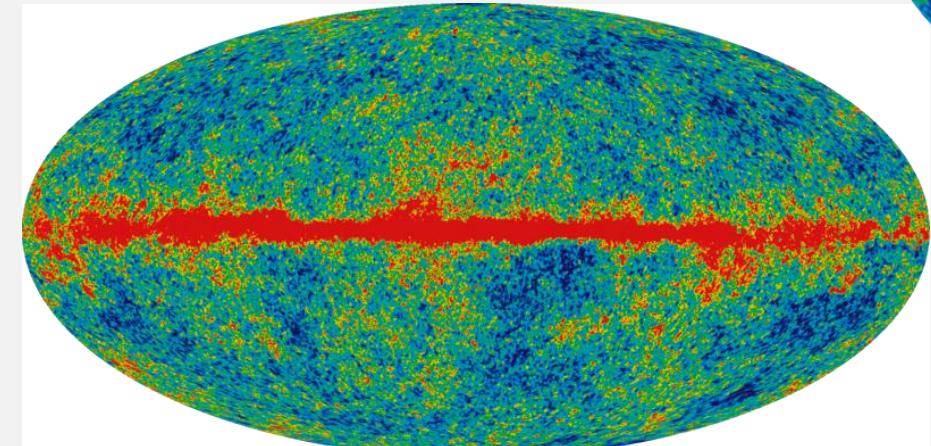


33 GHz with foreground

Final result, 7-years data
foreground subtracted

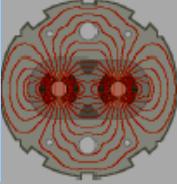


94 GHz with foreground



Inhomogeneities $\sim 10^{-4} - 10^{-5}$

WMAP results



◆ CMB power spectrum

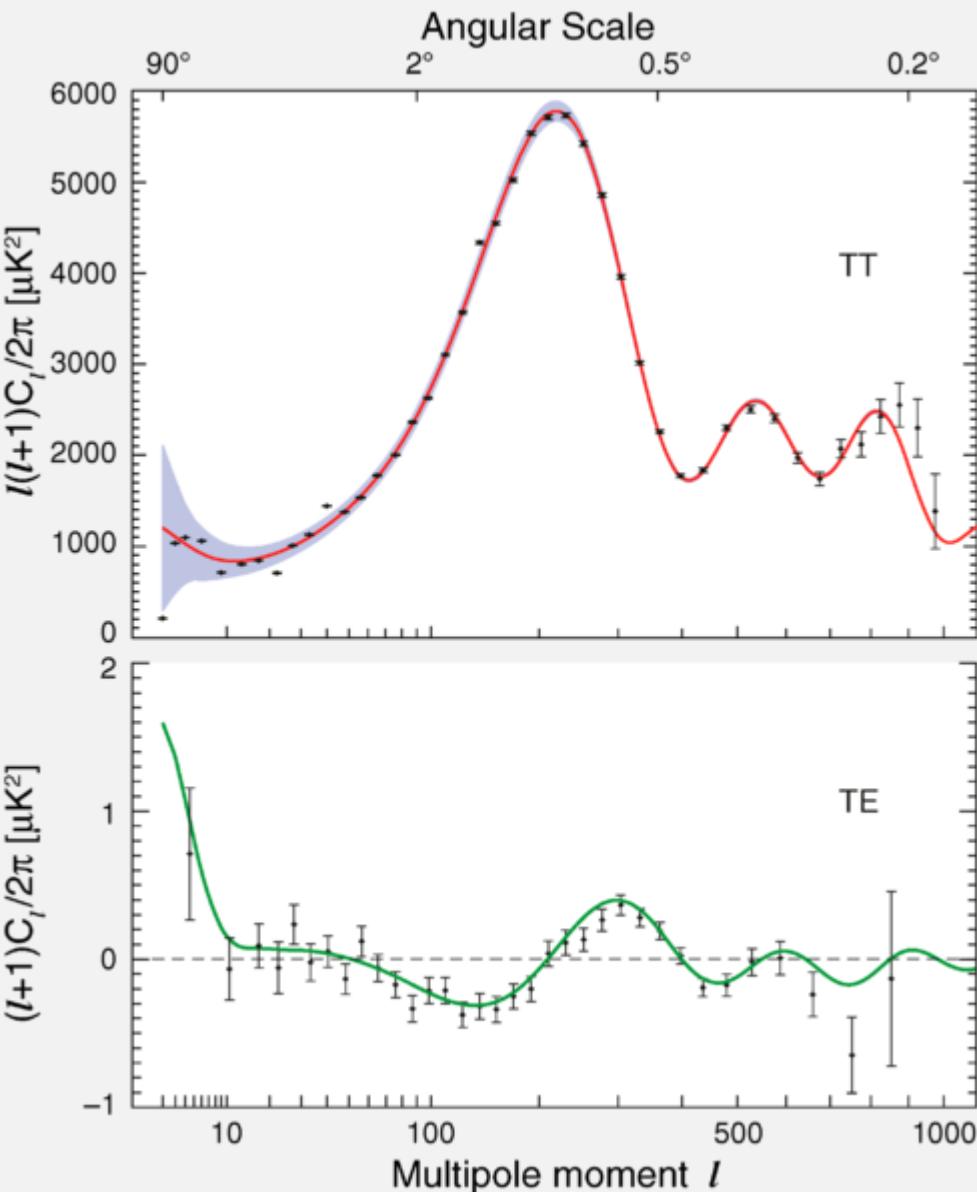
dipole moment gives our velocity 627 ± 22 km/s

age of the Universe
 $13.75 \pm 0.11 \times 10^9$ years

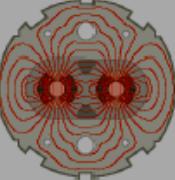
total density of the Universe
 $1.0023^{+0.0056}_{-0.0054}$

Hubble's constant
 $70.4^{+1.3}_{-1.4}$ (km/s) / Mpc

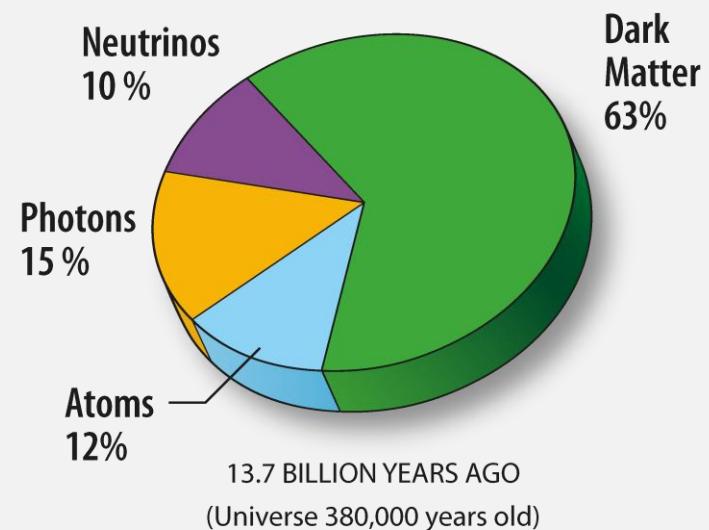
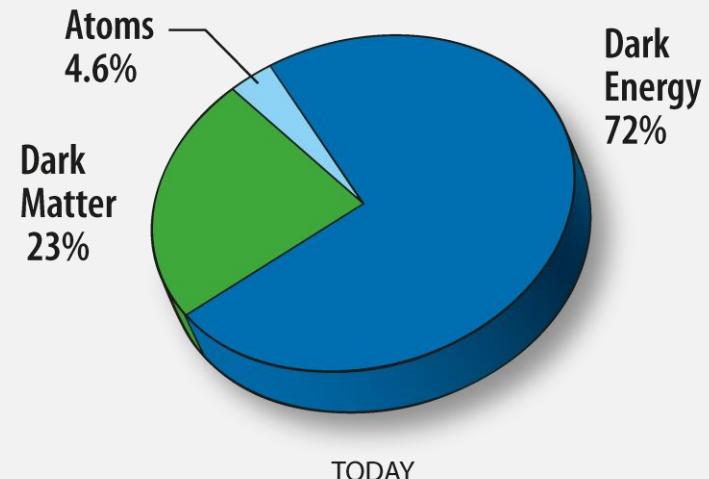
CMB temperature 2.725 K



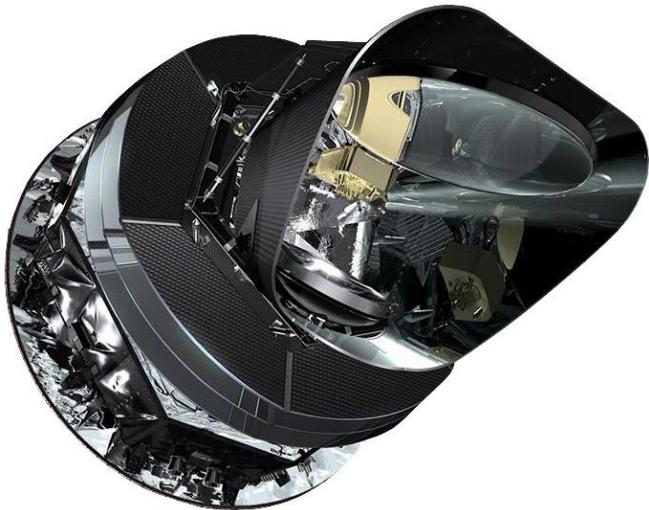
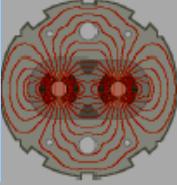
Composition of matter



- ◆ today 72% of matter of the Universe – dark energy
- ◆ before $\sim 7 \times 10^9$ years the Universe accelerated its expansion
- ◆ vacuum energy? scalar field? cosmological constant?
- ◆ 23% is (cold) dark matter, what is it?



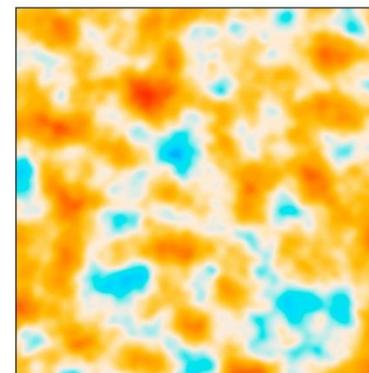
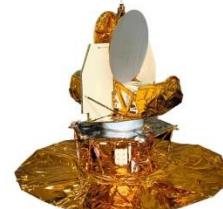
Planck satellite



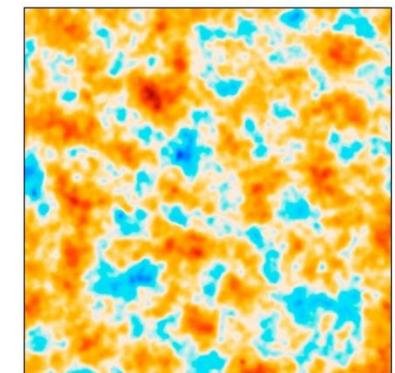
Porovnanie rozlíšania:
10 x lepšie rozlíšenie než WMAP
9 frequency band (WMAP 5)



COBE

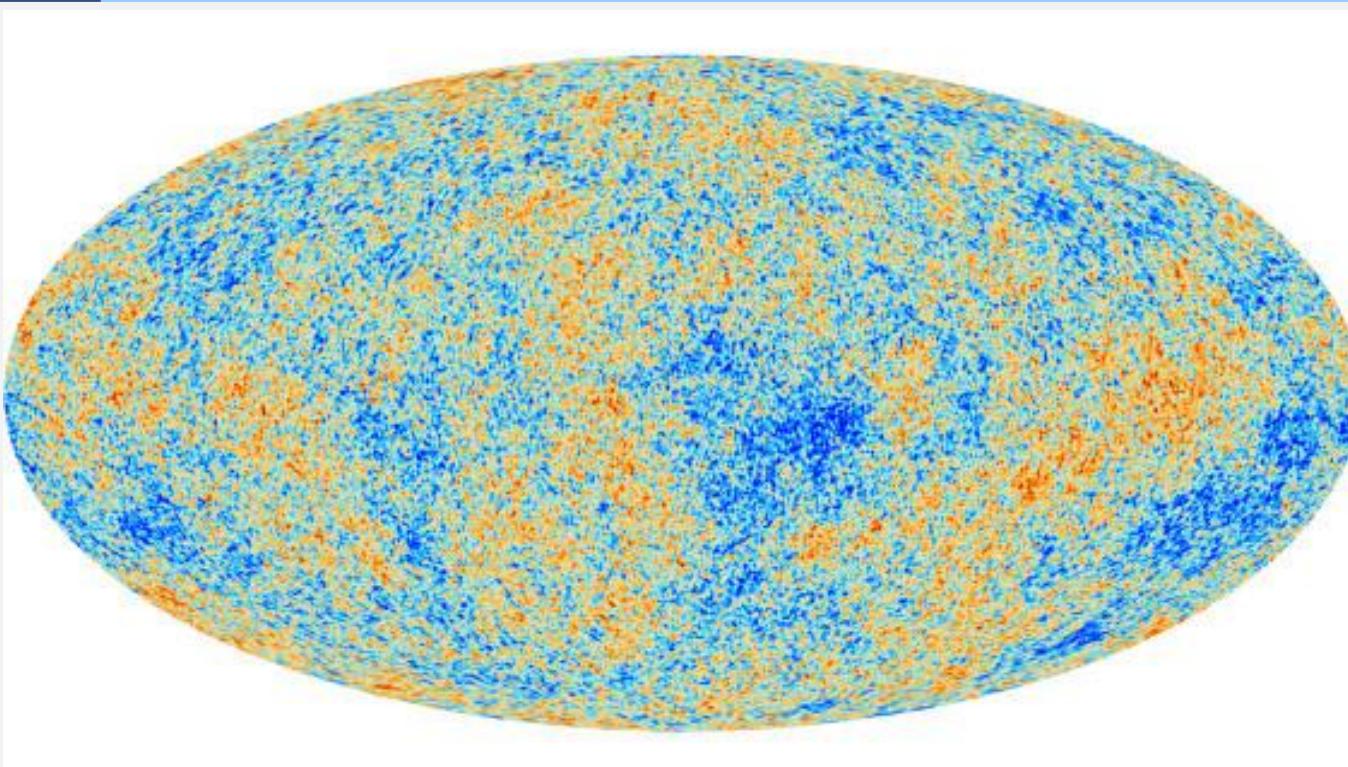
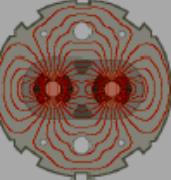


WMAP



Planck

Planck – prvé výsledky



Age of Universe: $13.798 \pm 0.037 \times 10^9$ years

Hubble constant: 67.80 ± 0.77 (km/s)/Mpc

Ordinary matter: 4.9 %

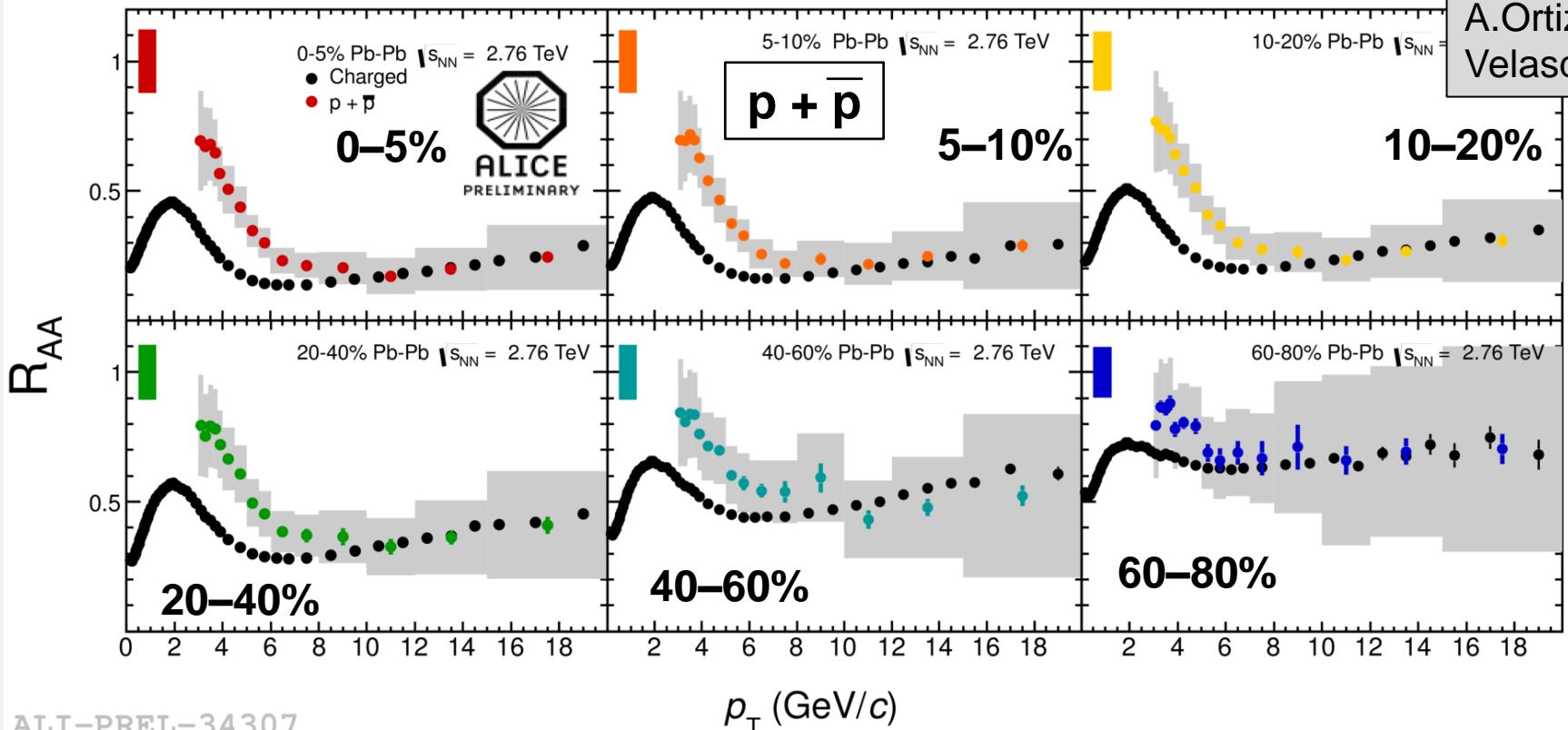
Dark matter: 26.8 %

Dark energy: 68.3%

Identified particles at intermediate p_T

- charged particles
- ● ● ● ● different centralities for identified particles

Talks by
M.Ivanov
A.Ortiz
Velasquez

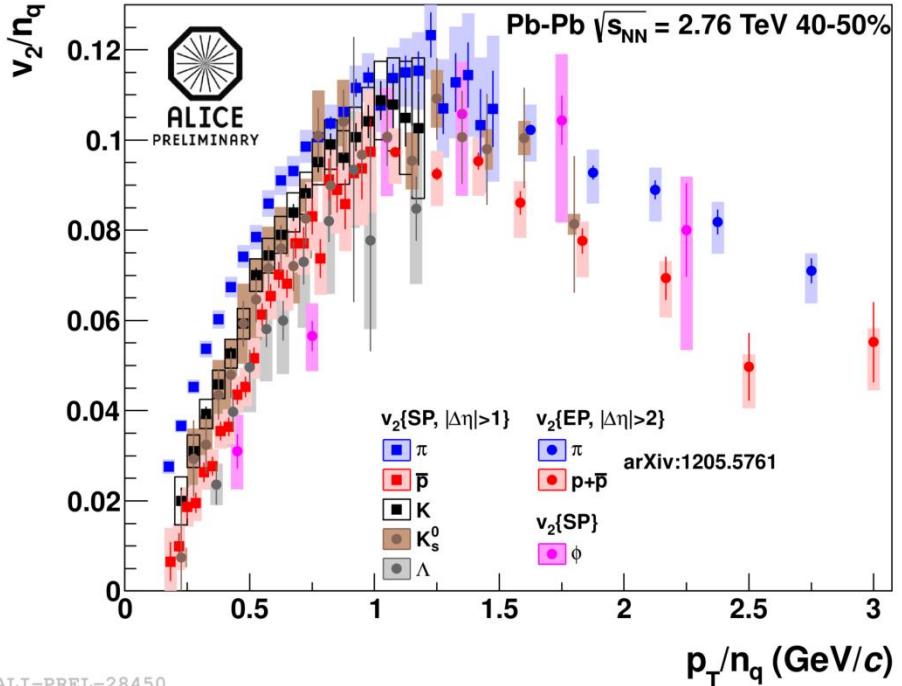
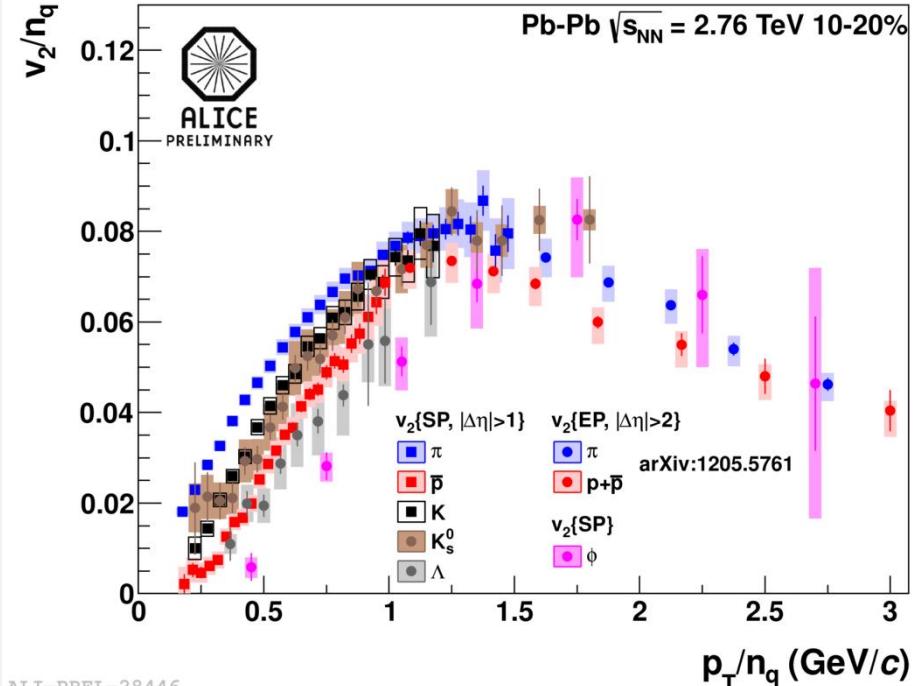


ALI-PREL-34307

For p_T below ~ 7 GeV/c: $R_{AA}(\pi) < R_{AA}(h^\pm)$, $R_{AA}(K) \approx R_{AA}(h^\pm)$, $R_{AA}(p) > R_{AA}(h^\pm)$

At higher p_T : R_{AA} are compatible

Identified-particle v_2



ALI-PREL-28446

ALI-PREL-28450

v_2 for π , p , K^\pm , K_s^0 , Λ , ϕ (not shown for Ξ , Ω)
 ϕ at low p_T ($< 3 \text{ GeV}/c$) follows mass hierarchy
– at higher p_T joins mesons
overall qualitative agreement with hydro up to
 p_T 1.5–3 GeV/c (π – p); quantitative precision
needs improvements – hadronic afterburner

$v_2\{\text{SP}, |\Delta\eta|>1\}$ $v_2\{\text{EP}, |\Delta\eta|>2\}$

π \bar{p} K K_s^0 Λ

$p+\bar{p}$

$v_2\{\text{SP}\}$ ϕ

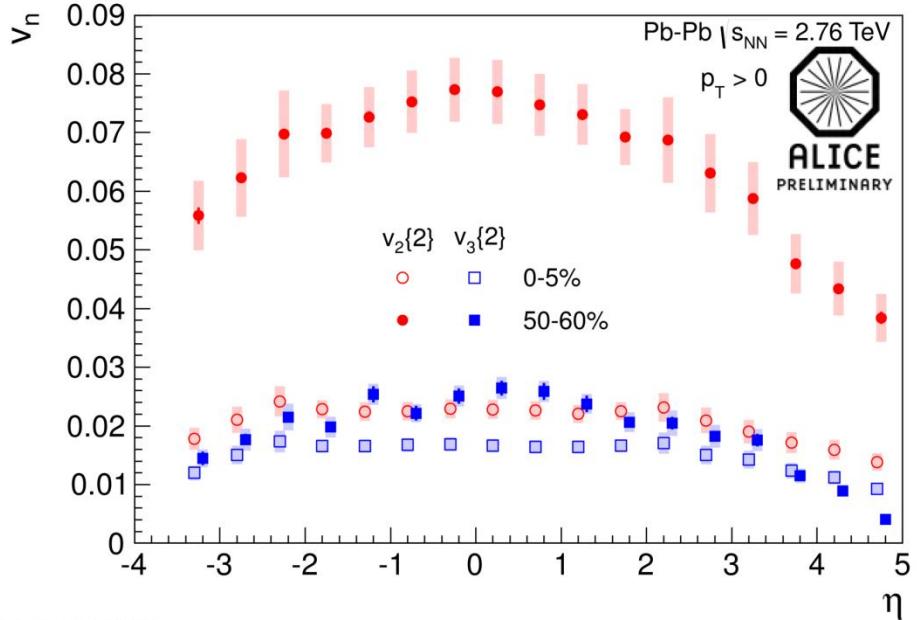
arXiv:1205.5761

Talks by
S.Voloshin
F.Noferini

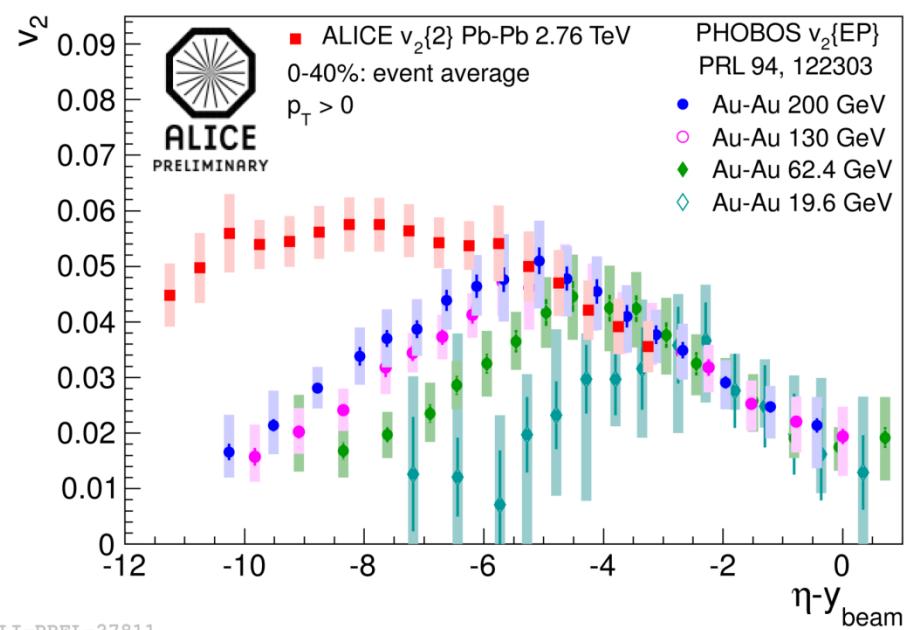
$n_q(m_T)$ -scaling worse than at RHIC

$n_q(p_T)$ -scaling at $p_T > 1.2 \text{ GeV}/c$ violation 10–20%

v_2 and v_3 versus η



ALI-PREL-28033



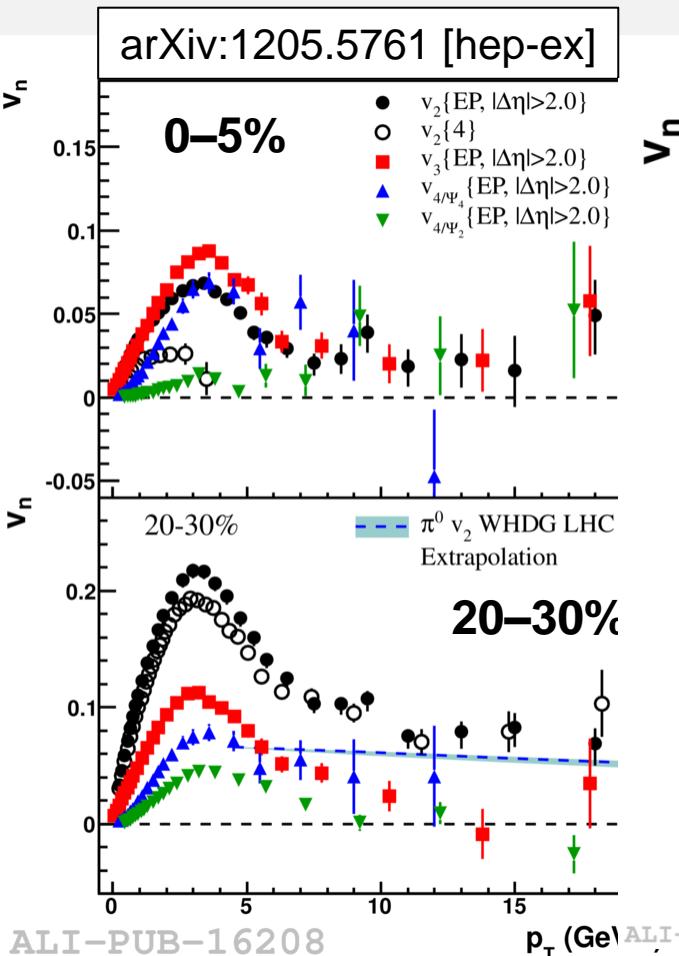
ALI-PREL-27811

v_2 and v_3 measurements extended up to $\eta = 5$
observed plateau in pseudorapidity ($|\eta| < 2$)
very good agreement between ALICE and CMS for v_2 in $|\eta| < 2.4$

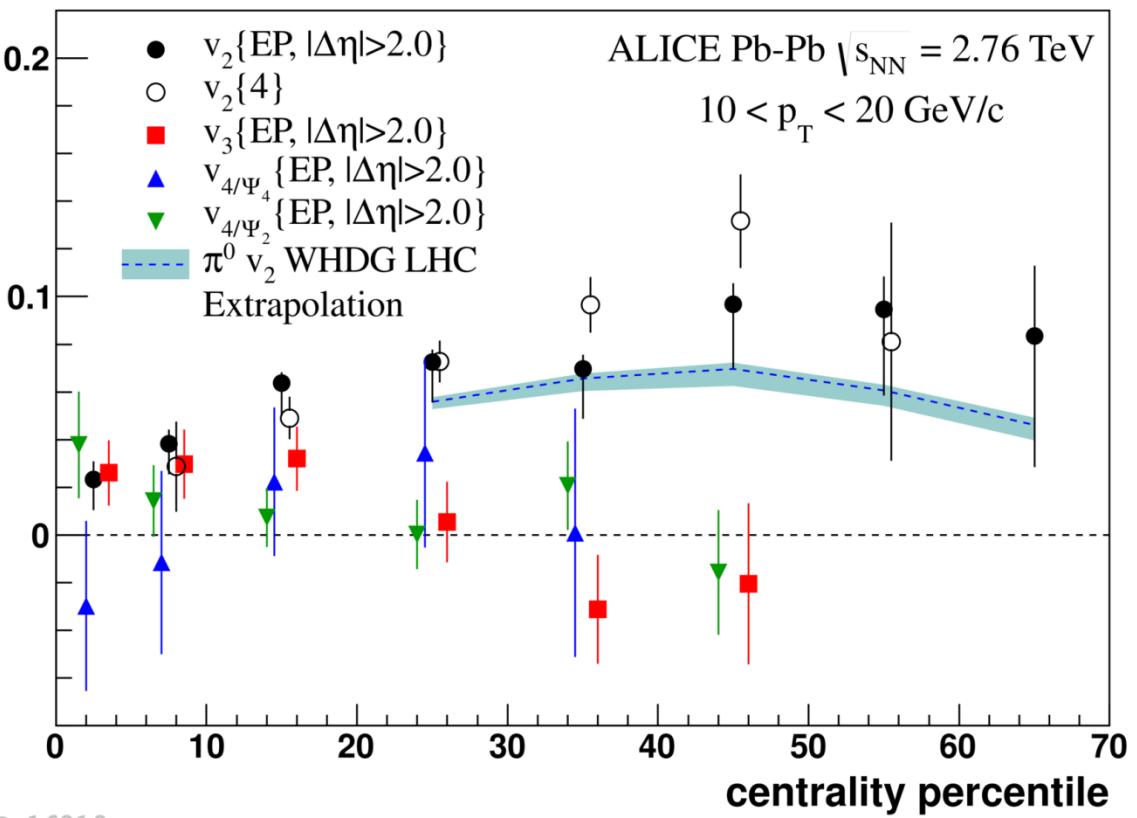
consistent with longitudinal scaling in $\eta - y_{\text{beam}}$ with PHOBOS data

Talks by
S.Voloshin
A.Hansen

v_2, v_3, v_4 versus p_T



----- W.Horowitz,M.Gyulassy, J.Phys. G38 124114

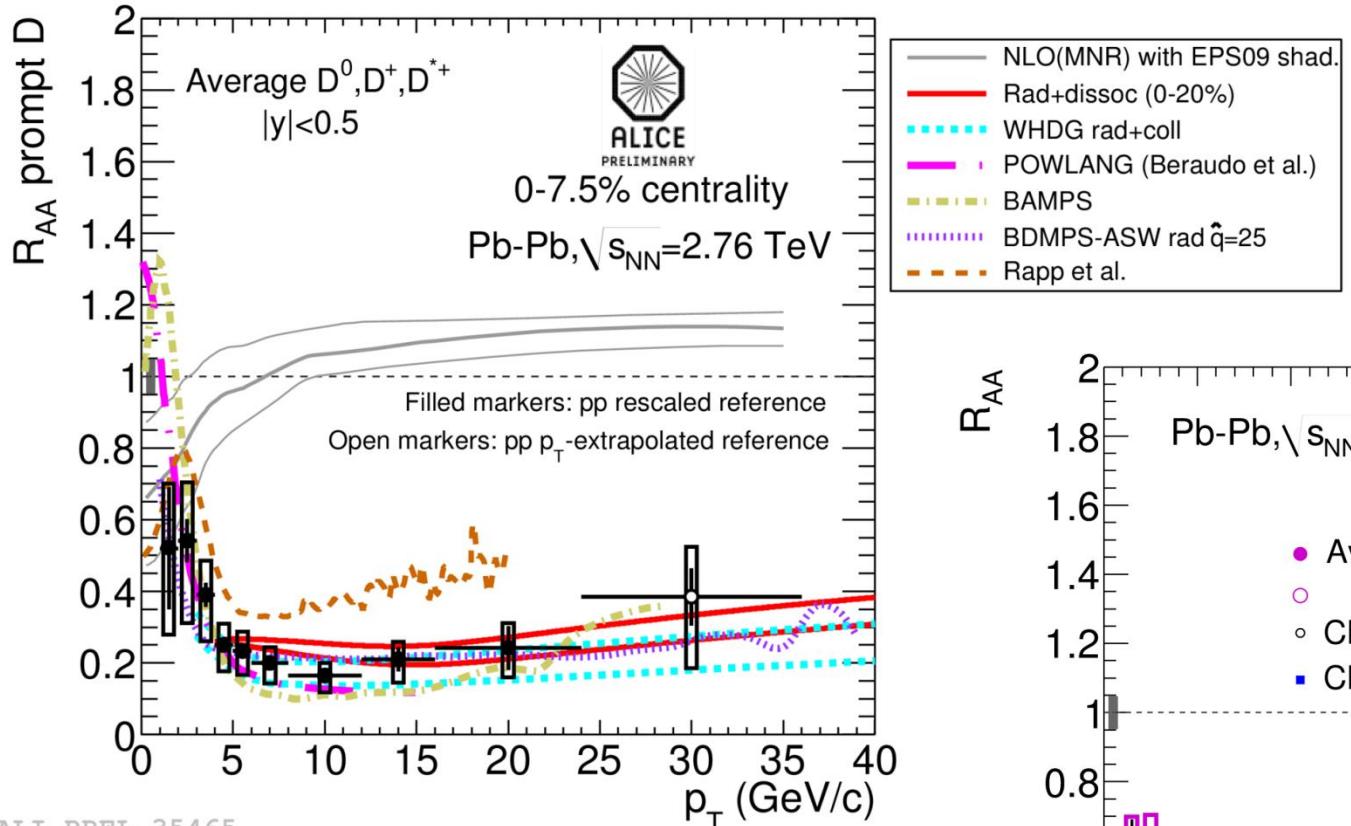


v_n measurements up to 20 GeV/c – where dominated by jet quenching
 Non-flow effects suppressed by rapidity gap or using higher cumulants
 Non-zero value of v_2 at high p_T both for $\Delta\eta > 2$ and 4-particle cumulant

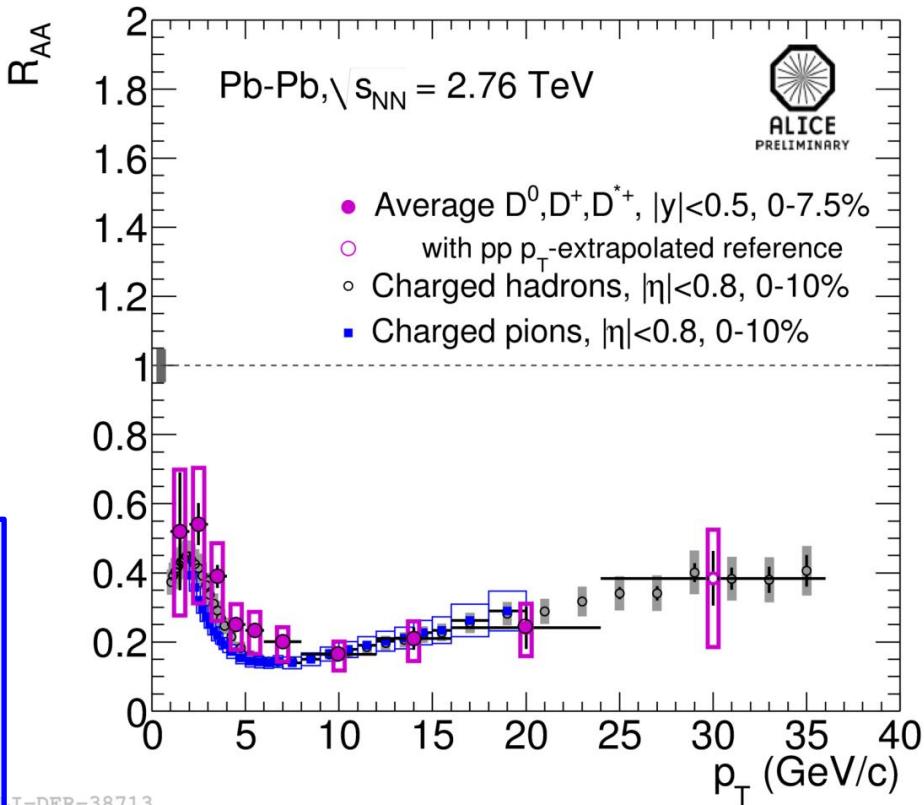
Talk by
S.Voloshin

v_3 and v_4 diminish above 10 GeV/c – indication of disappearance of fluctuations at high p_T

D meson R_{AA}



Talks by
Z.Conesa del Valle
A.Grelli

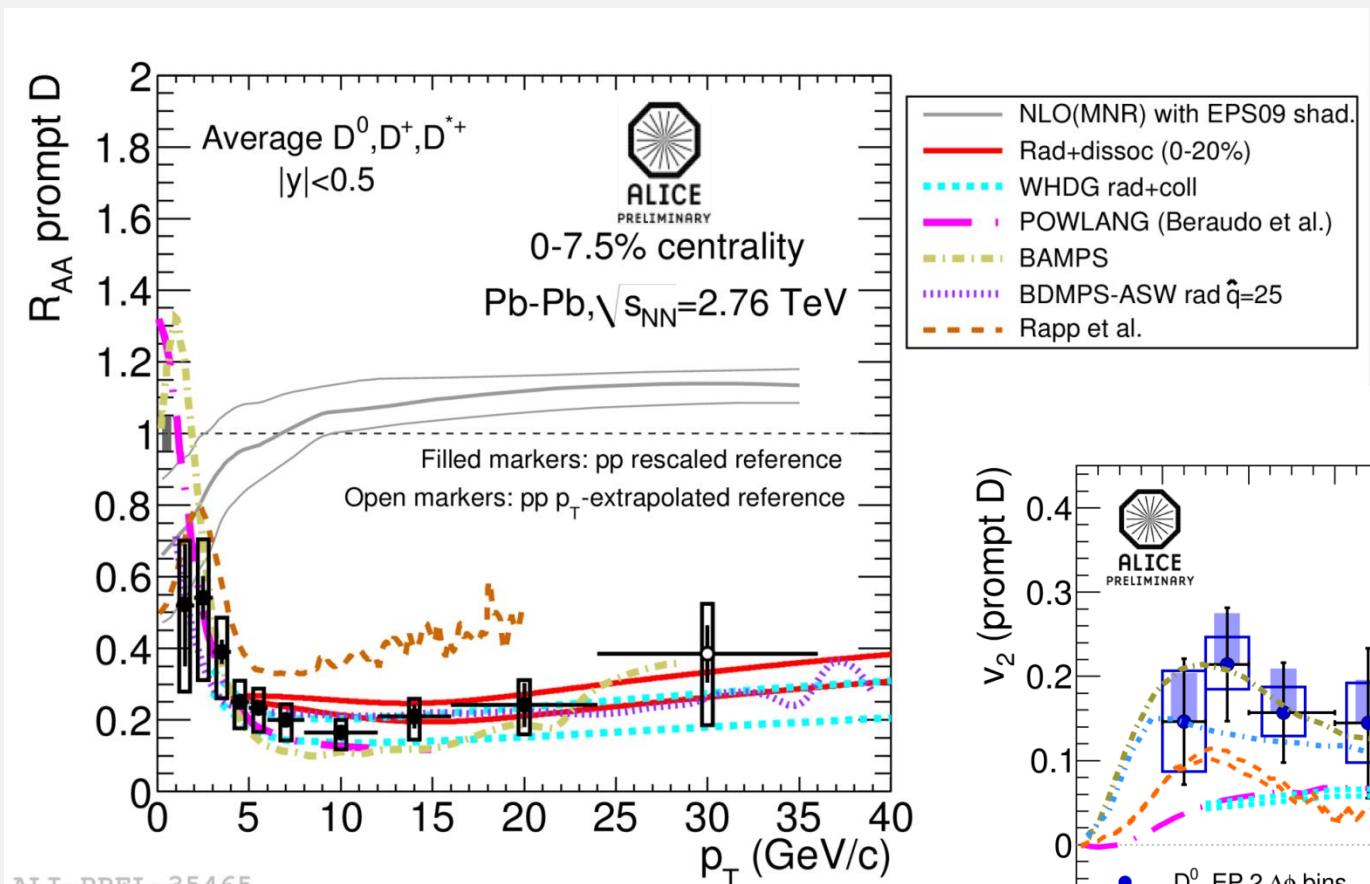


Average D-meson R_{AA} :

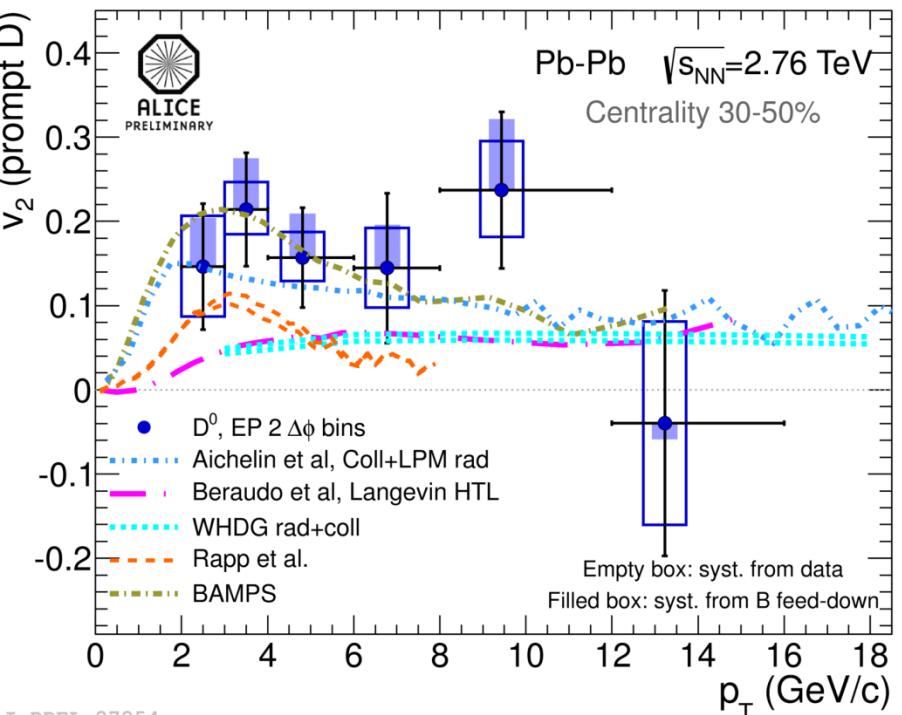
- $- p_T < 8$ GeV/c hint of slightly less suppression than for light hadrons
- $- p_T > 8$ GeV/c both (all) very similar no indication of colour charge dependence

ALI-PREL-35465

D meson v_2



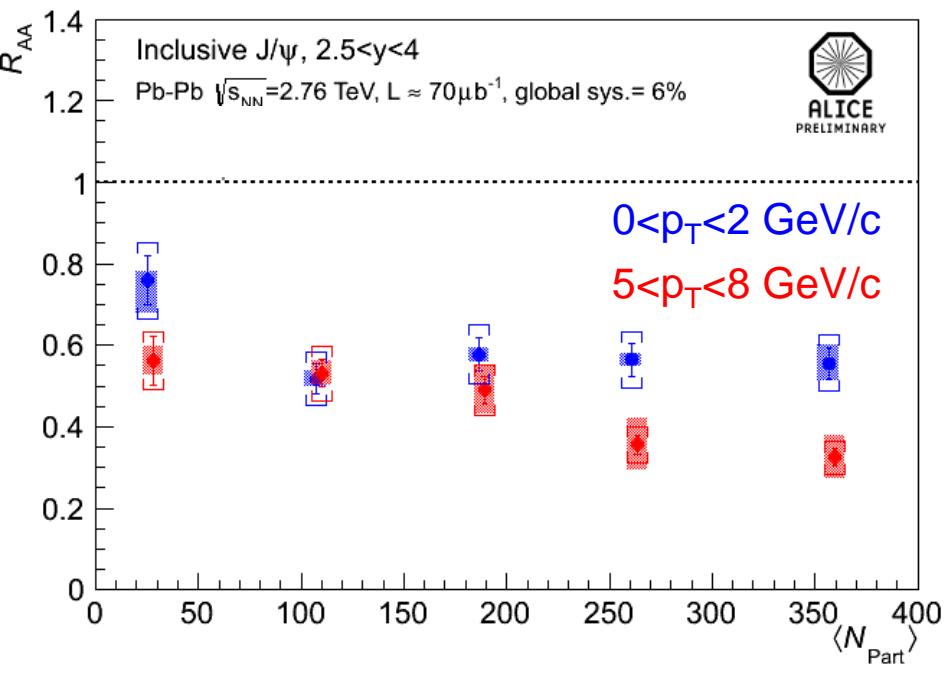
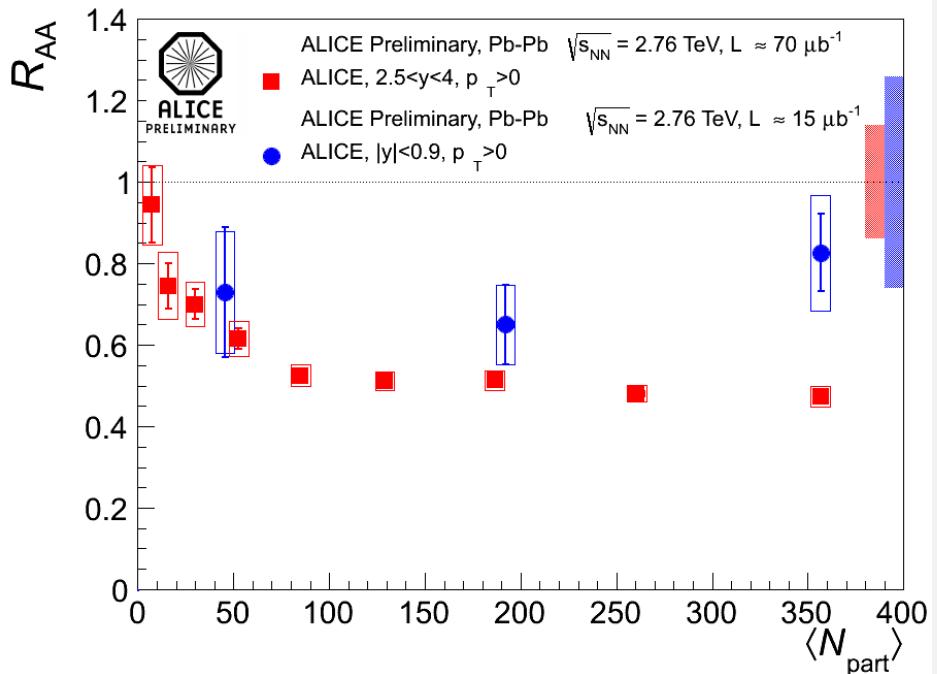
Talks by
Z.Conesa del Valle
D.Caffarri



Non-zero D meson v_2 observed
Comparable to that of light hadrons
Expressed as event-plane dependent R_{AA}

Simultaneous description of R_{AA} and v_2
c-quark transport coefficient in medium

J/ ψ R_{AA} centrality dependence



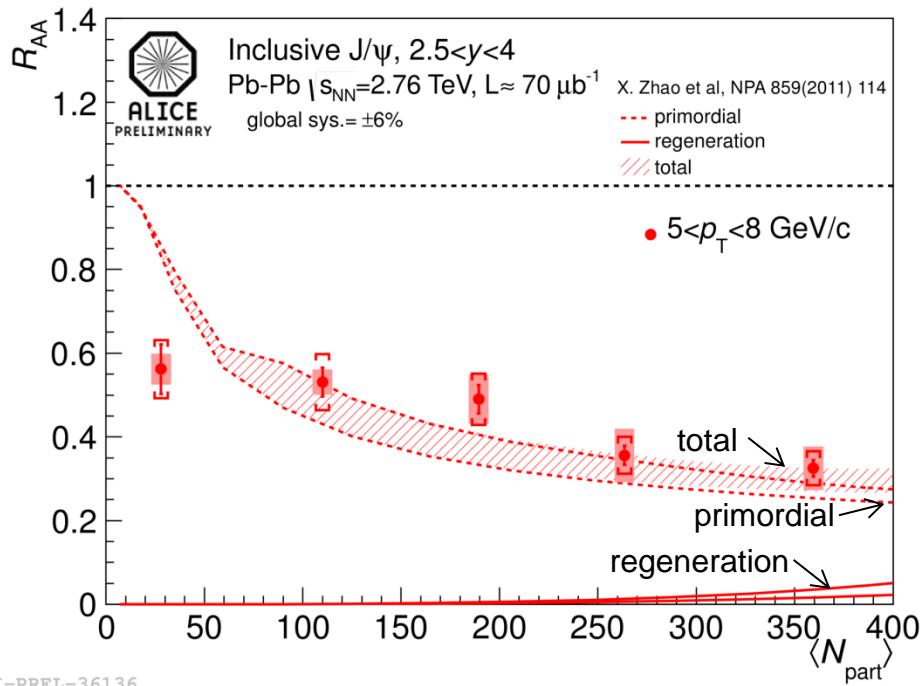
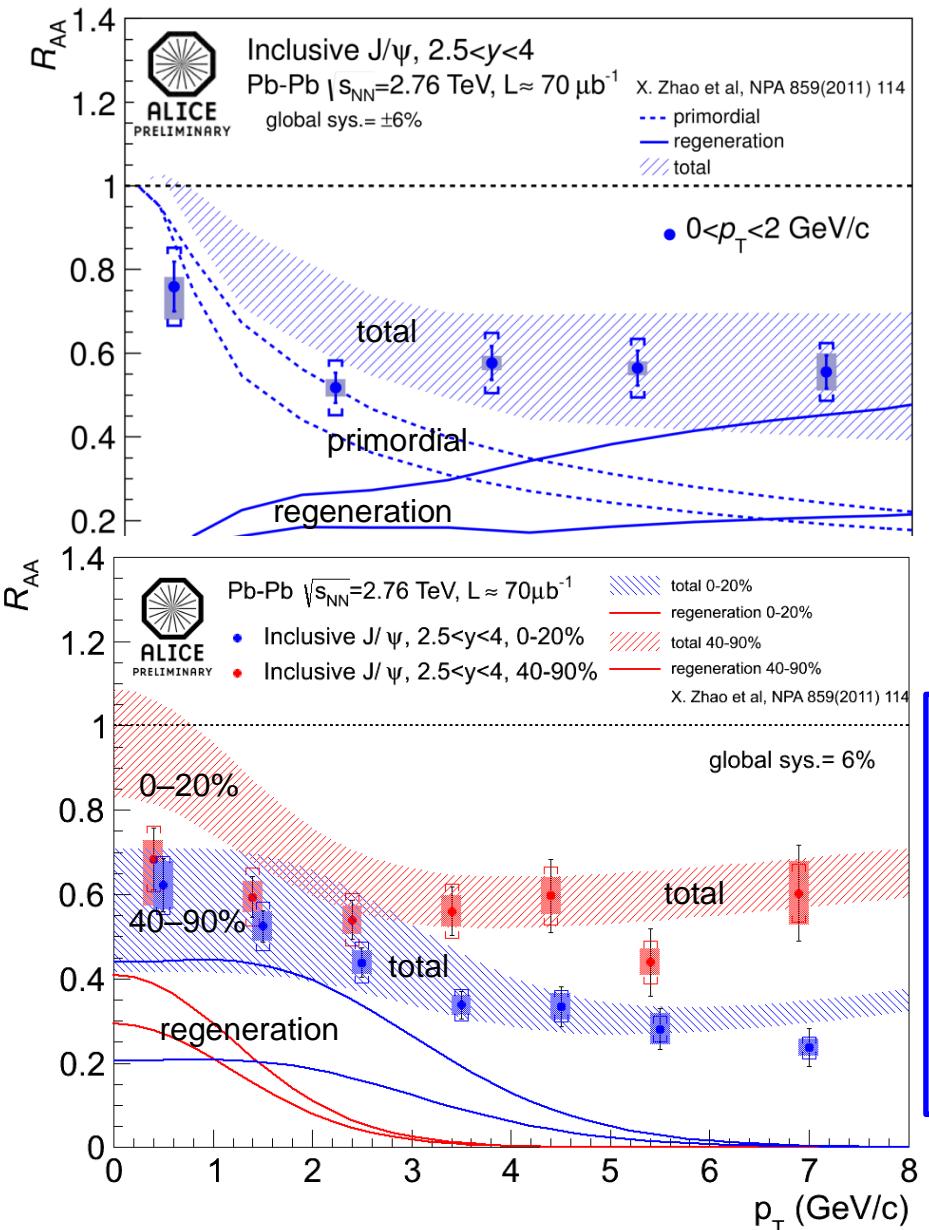
J/ ψ suppression measurements both in central and forward regions

- from $N_{\text{part}} > 100$ suppression independent of centrality
- in central collisions, less suppression than at RHIC
- at low p_T (< 2 GeV/c) less suppression than at high p_T , especially in more central collisions

Indication of J/ ψ regeneration at low p_T ?

Talks by
E.Scomparin
R.Arnaldi
I.Ch.Arsene

J/ ψ R_{AA} vs. centrality and p_T

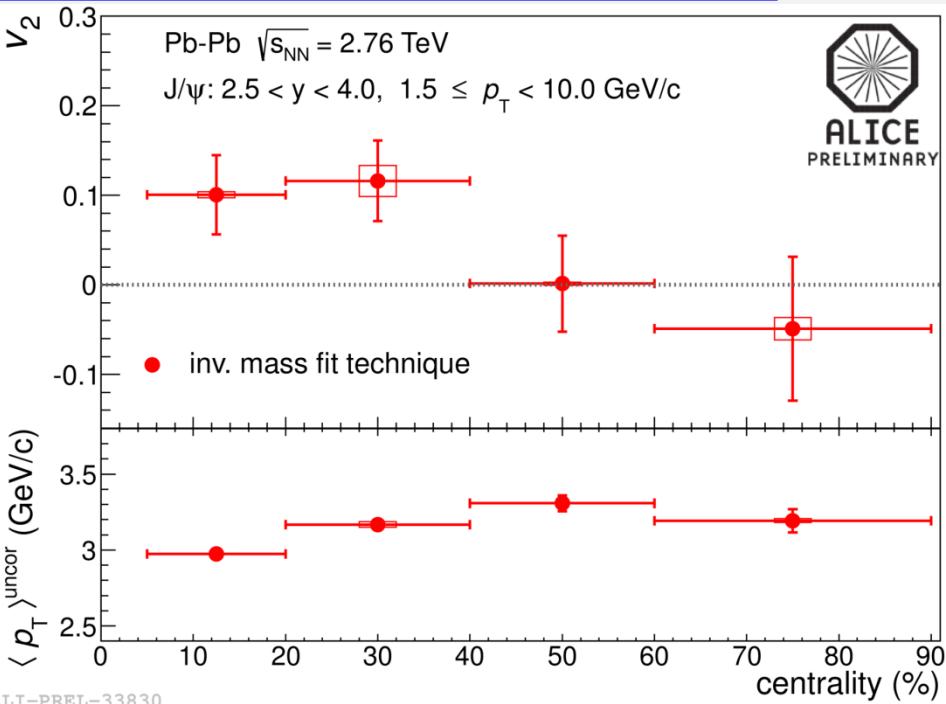
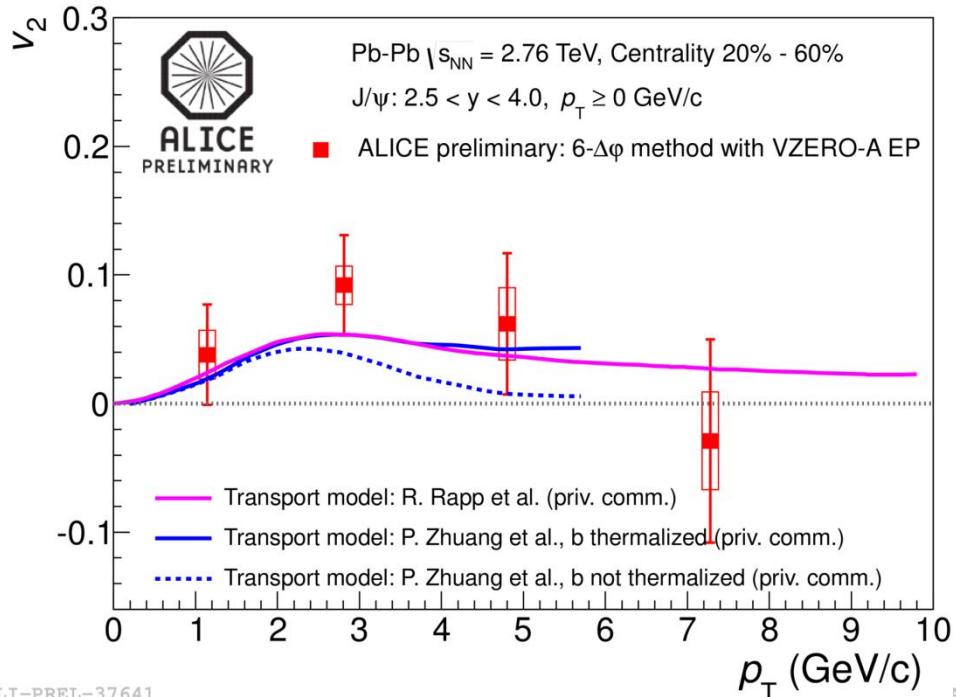


Comparison to regeneration model:
X.Zhao, R.Rapp NPA 859 114

Different suppression pattern at low/high- p_T
At low p_T ~50% J/ ψ from recombination
Fair agreement for different centralities
Statistical hadronization model also
describes the data: P.Braun-Munzinger et al.

Talks by E.Scomparin, R.Arnaldi

J/ ψ elliptic flow

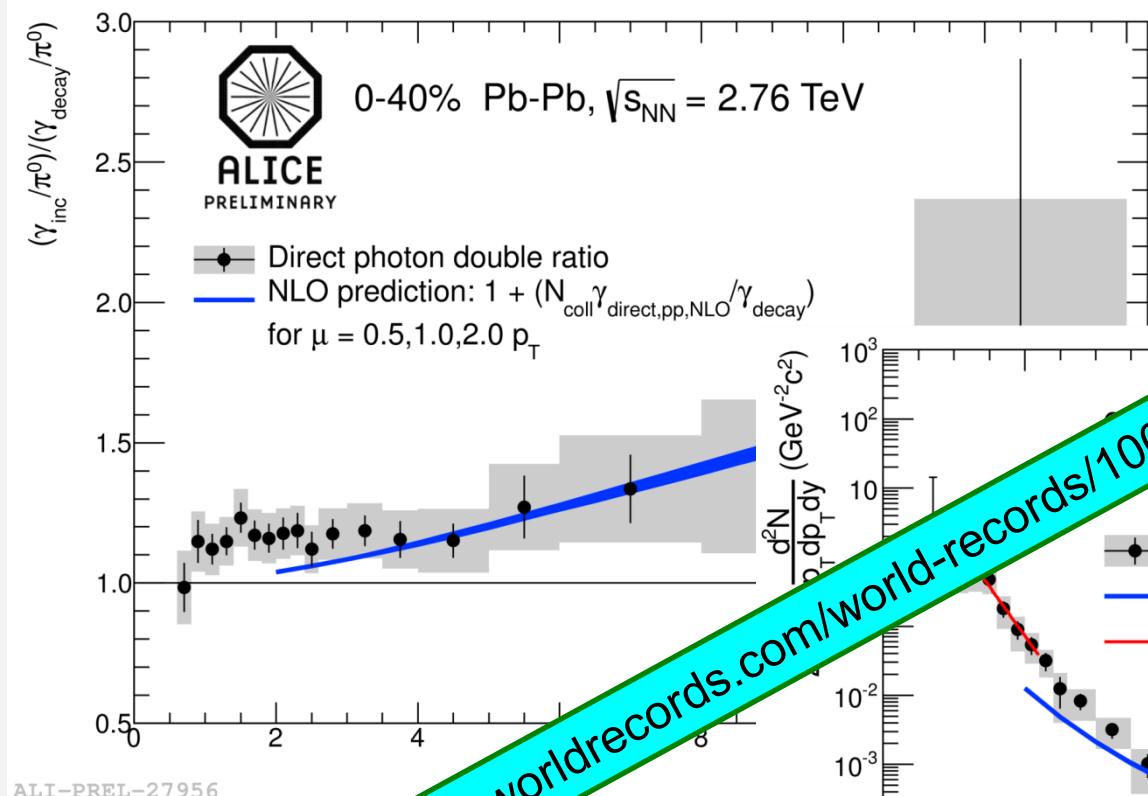


J/ ψ produced by recombination of thermalized c-quarks should have non-zero elliptic flow

- measurements give a hint for non-zero v_2
- qualitative agreement with transport models, including regeneration
- complementary to indications obtained from J/ ψ R_{AA} studies

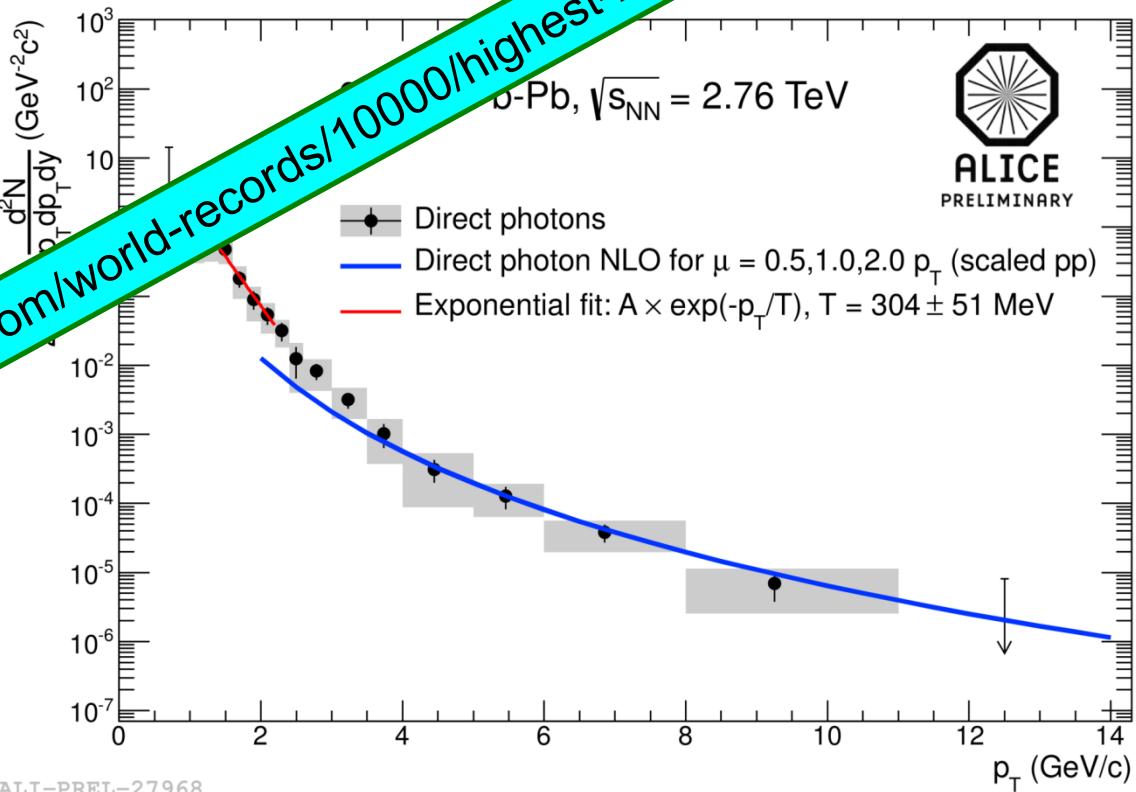
Talks by
E.Scomparin
H.Yang

Direct photon production



$p_T < 2$ GeV/c
~20% of direct photons
 $p_T < 2$ GeV/c
agreement with N_{coll} -scaled NLO

Exponential fit for $p_T < 2.2$ GeV/c
inv. slope $T = 304 \pm 51$ MeV
for 0–40% Pb–Pb at $\sqrt{s} = 2.76$ TeV
PHENIX: $T = 290 \pm 19$ MeV
for 0–20% Au–Au at $\sqrt{s} = 200$ GeV



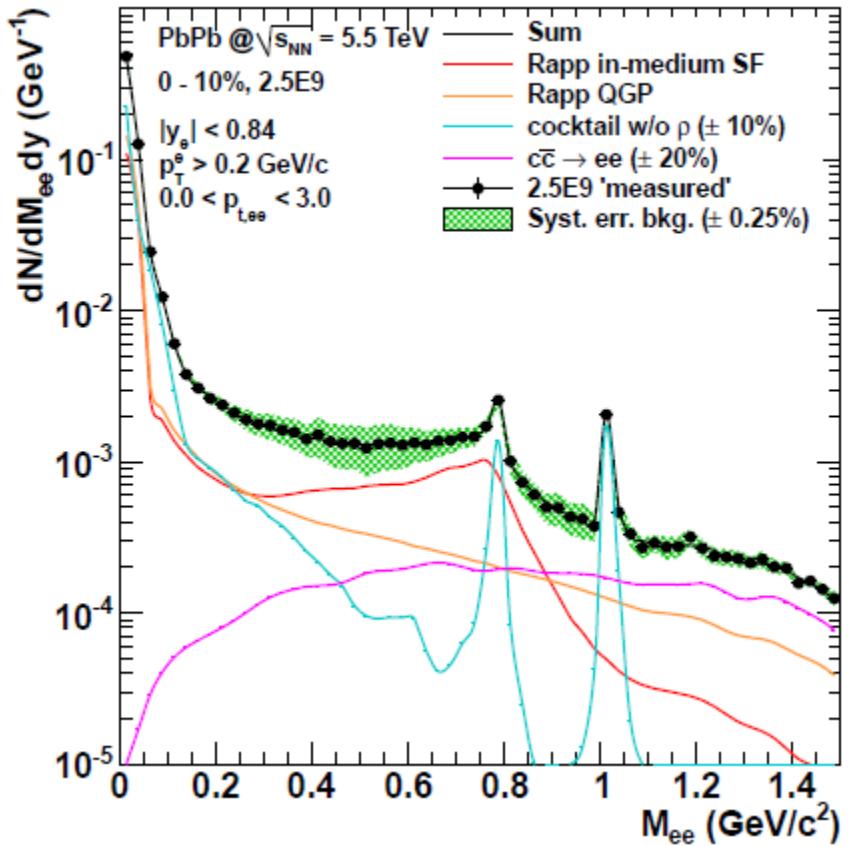
ALICE future plans

Precision measurement of the QGP parameters at $\mu_b = 0$ to fully exploit scientific potential of the LHC – unique in:

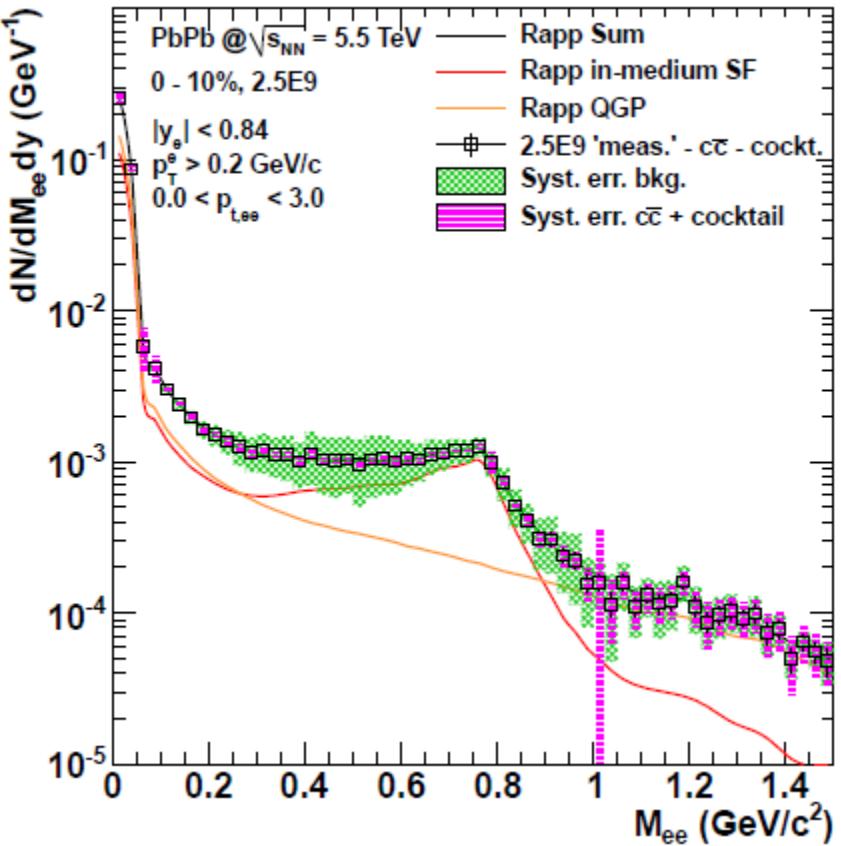
- large cross sections for hard probes
 - high initial temperature
-
- Main physics topics, uniquely accessible with the ALICE detector:
 - measurement of heavy-flavour transport parameters:
 - study of QGP properties via transport coefficients (η/s , \hat{q})
 - measurement of low-mass and low- p_T di-leptons
 - study of chiral symmetry restoration
 - space-time evolution and equation of state of the QGP
 - J/ψ , ψ' , and χ_c states down to zero p_T in wide rapidity range
 - statistical hadronization versus dissociation/recombination

ALICE dielectrons

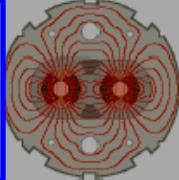
inclusive dielectron invariant mass



... excess after subtraction



new ITS and high-rate upgrade, with “tight” impact parameter cut...



● Common Questions

⇒ generation of mass

- ★ elementary particles => Higgs => ATLAS/CMS
- ★ composite particles => QGP => ALICE

⇒ missing symmetries

- ★ SuperSymmetry: matter <-> forces => ATLAS/CMS
- ★ Chiral Symmetry: mass of light quarks => ALICE
- ★ CP Symmetry: matter <-> antimatter => LHC-B

● Different Approaches

⇒ 'Concentrated Energy'

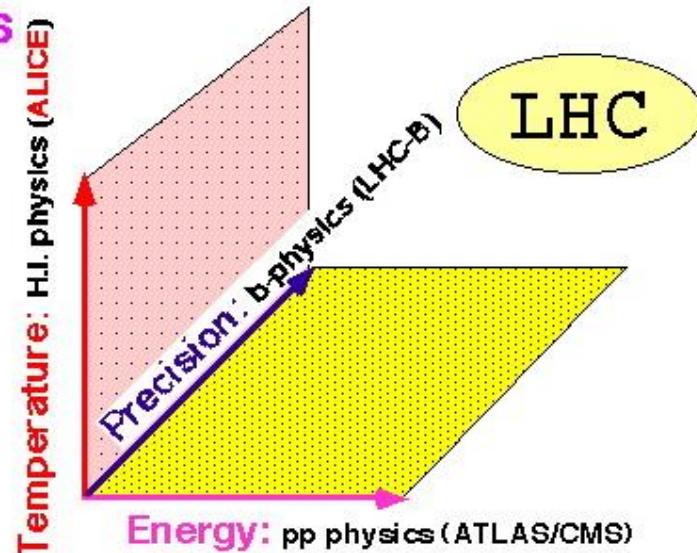
=> (single) high mass particles

⇒ 'Distributed Energy'

=> interaction between matter & vacuum

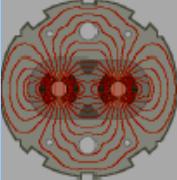
⇒ 'Borrowed Energy'

=> indirect effects of very high mass particles





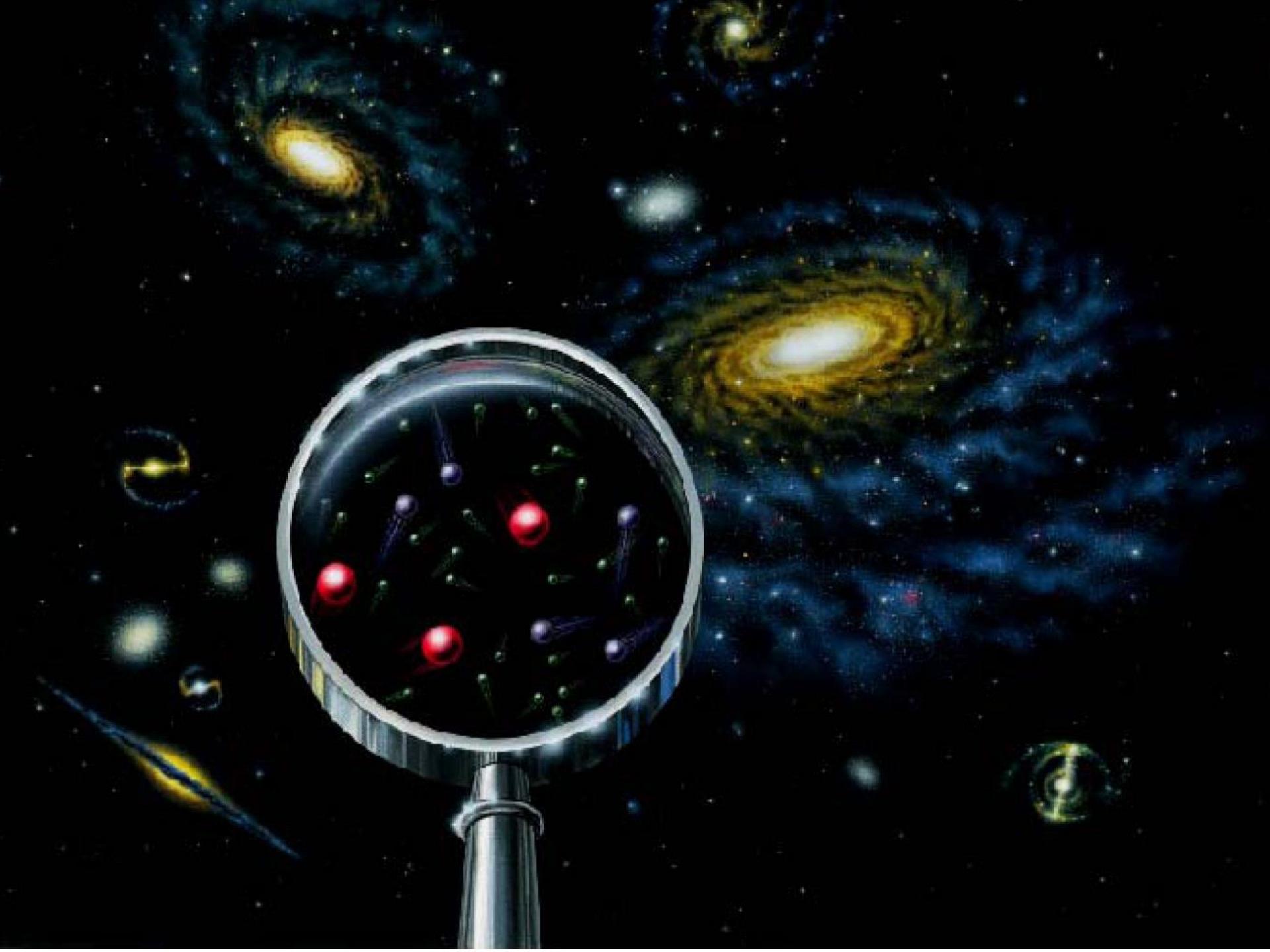
BANGS



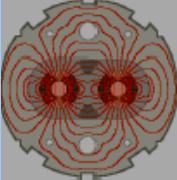
Big Bang \leftrightarrow Little Bangs

- The matter content of the Universe
 - Dark matter
 - Dark energy
 - Origin of matter
- Experiments at particle colliders
 - Early Universe
 - Supersymmetry
 - Matter-antimatter asymmetry

Learn particle physics from the Universe
Use particle physics to understand the Universe



Zhrnutie



- ◆ LHC a vsetky experimenty pracuju tri roky perfektne
 - intezity zvazkov a mnozstvo zaznamenanych zrazok ovela nad ocakavanie
 - prve fyzikalne vysledky publikovane
 - presne potvrdenie stardartneho modelu
 - Higgsov bozon najdeny – podla predpovedi hmotnost ~ 125 GeV
 - fluktacie v pociatocnom stave v zraskach tazkych ionov
 - dosiahnuta teplota v zraskach tazkych ionov ~ 300 MeV
- ◆ Looking forward to explore the ‘terra incognita’ at LHC



Hic sunt Leones !

Od veľkejho tresku k Leniu - sivočerné vesmíru v laboratóriu